COMPARATIVE STUDY OF STRUCTURAL BEHAVIOURS AND DESIGN OF LONG CYLINDRICAL CONCRETE SHELLS USING BEAM AND CLASSICAL METHODS

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in partial Fulfilment of the Requirements
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MASTER OF TECHNOLOGY

By PRABIR CHANDRA BASU

to the

DEPARTMENT OF CIVIL ENGINEERING

INDIAN INSTITUTE OF TECHNOLOGY KANPUR
SEPTEMBER 1975



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to my parents



CERTIFICATE

'COMPARATIVE STUDY OF STRUCTURAL BEHAVIOURS AND DESIGN OF LONG CYLINDRICAL CONCRETE SHELLS USING BEAM AND CLASSICAL METHODS', by Mr. Prabir Chandra Basu was carried out under my guidance and has not been submitted elsewhere for a degree.

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Dated: September, 1975

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NOTATIONS

Greek Letters

 α , α' = Roots of Characteristic eqns.

 β , β ' = Roots of characteristic eqns.

m = Const. for roots of characteristic eqns.

Const. for the solution of stress resultant

 ϕ_{k} = Semicentral angle

Ø = Angle at any circumpherial section

▼ = Differential opperator

 $\emptyset' = \emptyset_k - \emptyset$

Ux = Longitudinal stress

= Shear stress

 Ψ_1 to Ψ_2 = Trigonometrical and hyperbolic multiples

 θ = Rotation

δ = Ratio of the deformations

e = Error in 📆

Capital Letters

A = Area of cross section

A = Co-eff. matrix for arch analysis

 A_m = Const. for homogeneous solution

 B_1, B_2, B_3, B_4 = Shell const. for homogeneous solution

C = Constant vector for arch analysis

C_b = Beam constant or const. for perticular integral

E = Modulus of elasticity

```
P
                       Nodal forces
                  =
                       Coeff. of Torsional rigidity
       G
       H
                       Harizontal end reaction
        I
                       Moment of Inertia
        J
                       Polar moment of Inertia
K11, K21, K22, K31
                       Consts for arch analysis
        L
                       Span
       \mathbf{L}_{\mathbf{c}}
                       Chord length
       M_{\mathbf{m}}
                       End moment for shell or bending moment
       M
                       Transverse moment for shell or banding moment
                       Tersional moment for shell
    Mox My
        Mx
                       Longitudinal moment for shell
                       Axial thurst for arch
        Nø
                       Inplane force of shell
                       Principal force
       M, M
                       Longitudinal force for shell
    Nøx, Nxø
                       Shear force for shell
       Pe
                       Eff. prestressing force
                       First moment of area or const.for solutionof
        Q
                       characterestic equation of shell
       Q_{\mathbf{x}}
                       Radial shear
       Qø
                       Shear in the Ø direction
       R
                       Radins
                       Shear force
        Sp
                       End correction force. (Shear)
```

Const., Radius to thickness ratio

T

U = Strain energy

V = Vertical correction force

X,Y,Z = Co-ordinate system

Z = Distance of C.G.

 Z_v = Distance of fiber

Z_r = Distance of the C.G. of the element from the whole section C.G.

Small Letters:

a = Width of arch, const.of homogeneous solution of shell eqn.

b = Width of edgebeam, const. of homogeneous solution of shell eqn.

c = Coeff. of defermations, stress resultants, const. of homogeneous solutionof shell eqns, Consts.

d = depth. of edge beam, const. of homogeneous
 of shell eqn.

ds = Elementery arc length

dy = Horizontal projection of ds

dz = Vertical projection of ds,

e = Eccentricity of tendon for beam theory

e = End eccentricity of tendon for beam theory

en = End eccentricity of tendon in edge beam

e_m = Mid eccentricity of tendon in edge beam

eh = Horizontal eccentricity of the shell central line from the centre line of edge beam

e_v = Vertical eccentricity of centre line of shell from centre line of edge beam

g = Sag

h = Total depth of trasverse section of the shell with edge beam

```
Constants. n_{\pi}x/1, no. of slices of arch analysis
      k
    k_m, k_n
                    Const. for arch analysis
      1
                    Span
              =
      1<sub>c</sub>
                    Choid length
              霊
                    Root of the charecteristic eqn., nπ x/1,
      m
              =
                    total no. of nodes.
                    No. of fourier terms, k-1
      10
              =
                   Nondimenerionalised coeff. for prestressing force,
      p
                    external loading
                    Intensity of loading on shell
               =
      q
                    Thickness of shell
      t
               =
                    Displacements
      u
               =
                    Deformations
      W
               =
Subscripts
                    For N when used in first place, for V when used
      1
               *
                    in second place
      2
                    For N when used in first place, for S when
               =
                    used in second place
                    For Nx when used in first place, for H when
      3
                    used in second place
      4
                    For Nx when used in first place, for M when
                    used in second place
      b
                    Edge beam
               ≢.
                    Shell without edge beam
      C
                    Effective
      е
      ſ
                    force
```

Nodal point, V-

i

==

g = Load on the surface area of shell

h = Horizontal direction, due to H

i = Nodal points,

k = Shell edge

1 = Line load

m = Due to rotation, membrane action

r = Radial direction

s = Shear force, shell with edge beam

u = Uniformly distributed load on chord length

v = Vertical direction

w = Due to deformation

x,y,z = In the direction of co-ordinate axes

I = Const. for moment of inertia

H = Due to horizontal force

M = Due to moment

V = Due to vertical force

Z = Const. for C.G.

 \emptyset = In the direction of \emptyset

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SYNOPSIS

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The thesis presents the comparative study of the structural behaviour and design of long cylindrical concrete shell having span to radius ratio greater than 3, are assumed as long shells and such shells with edge beam are considered for the investigation. Both reinforced and prestressed concrete shells are studied. Only edge beams of prestressed shells were provided with prestressing force. The response of the stress resultants and deformations are taken as structural behaviour parameters.

Paramatric study was made to investigate the influence of different shell parameters and prestressing on the behaviour evaluated by both the theories and to find out the domain of applicability of the beam theory. It was found that in case of reinforced concrete shell, beam theory is applicable for all

normal values of radius to thickness, depth of edge beam to span ratios, semicentral angle and span to radius ratio greater than 3., within 4 to 15 percent of error.

Prestressing reduces the stress resultants and deformations with a suitable combination of depth of edge beam, and wideccentricities of the tendon. The discrepency in the assumptions of beam theory increases with the increase of depth of edge beam and prestressing force. Restraints against the horizontal and rotational deformations yield the stress resultant which are in nature to compensate each other. No shell with coefficient of prestressing more or equal to 0.4 and depth of edge beam to span ratio more or equal to 0.04 could be designed by beam theory.

that

It was found the arch analysis by column analogy method considering a thin strip of the shell with edge beam is not efficient. A simplified formulation for arch analysis based on strain energy method, satisfying the compatibility of deformation at the interface of the shell and edge beam, is developed. With the help of this arch analysis a systematic and fairly simple method, detailed in six steps, is presented. The proposed method gives good results for single barrel shell.

A series of table is given to make the computational work easy for the proposed method.

CHAPTER-I

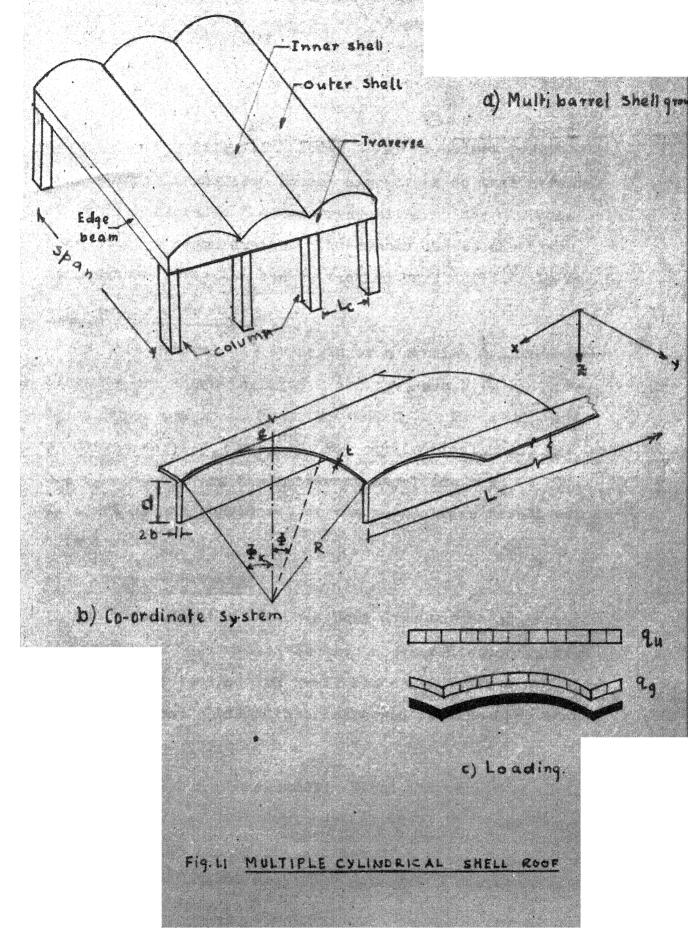
INTRODUCTION

1.1 GENERAL:

Shells are one of the economical, effective and elegant structural system for the purpose of roofing of a large unobstructed area. They tend to carry loads primarily by direct stress acting in their plane, because of their geometry and small flexural rigidity of skin which catagorise them to stressed skin structures. Among the different types of shell, cyllindrical shell is most popular for roofing purpose. They have extensive application in constructing airport hangers, industrial buildings, auditoriums, exhibition stalls, godowns, markets, garage units, bus terminus etc.

For small span, shells are: not provided with edge beam, but as span increases edge beams are provided to reduce the stresses. For very long span prestressing of edge beam yields the optimal utilization of structural properties of cylindrical shell roofs. Prestressing of edge beam has certain advantages as it reduces stress resultant and hence offers greater safety against buckling; deflections; consumptions of steel and dead weight of the structures.

With the suitable selection of end and mid eccentricity and



magnitude of prestressing force it is sometimes possible to eleminate transverse moment altogether so that stresses in shell reduces to those corresponding to the membrane plus a constant; compression. The cracks due to shrinkage, temperature variation can be tackled very efficiently by it.

1.2 STRESS RESULTANTS:

In general the stresses of a continuous structural elements are shown in Fig. 12 and they are N_X , $N_{X}\emptyset$, $N_{0}X$, $M_{X}\emptyset$, and $M_{X}\emptyset$ are considered as the main stress resultants in this investigation.

1.3 LITERATURE REVIEW:

Considerable work had been done on the analysis and design of Cylindrical Shell. However discussion here is restricted to the work available on analysis and design of reglindrical shell with prestressed edge beam by classical and 'Beam Theory'.

In 1949 H. Lundgren(3) first introduced a simple theory to analyse long explinarical shell by 'Beam Theory'. He compared the transverse section of a cylindrical shell as a curved section of a beam and analyse the shell for $N_{\rm x}$ and

 $N_{\mathbf{x}}$ considering the shell as a beam spanning from one sage transcribe to another in longitudinal direction. For the calculation of Mo and No he performed the arch analysis considering a unit width strip of shell as an arch loaded by the external loading and NxO which is calculated by beam calculation. The exch calculation that he did was for the interior shell and analysed it by column analogy method. J. Chinn (12) in 1959 demonstrated that the interior shell of a group of shells without edge beam even with $L/R \gg 1.11$ can be analysed by beam theory. For arch calculation he applied the elastic centre method. He also compared stress resultants obtained by beam theory and ASCE Manual (31). In the same year A.L Parme and H.W. Conner (14) introduced tables for the correction of stress resultants obtained by beam theory. Ramaswami (7) described almost all of the so called exact theory to analyse cylindrical shell. He also demonstrated the procedure of designing the explindrical shell with prestressed edge beam and analysed aninterior shell of a group of shell without edge beam having L/R=2 using column analogy method for arch calculation. Gibson(1) in 1968 presented the method to analyse the cyllindrical shell by assuming it as shallow shell. A detailed method to analyse a shell with prestressed edge beam was put by him and a companison of Nx obtained by exact method and beam

theory for shell without edge beam and with prestressed and non prestressed edge beam had been done by him. No. 31 of ASCE on the Design of Shell Roofs(27) contains a series of comprehensive tables and charts which reduces the lot of computational hazard. P. Dayaratnam (8) illustrated completely the procedure to design a single barrel shell with prestressed cage beam using ASCE Manual No. 31. D.P. Billington(2) discussed the classical shell theory by theory of shallow shell and use of ASCE Manual-31 simplified the design problem. He also demonstrated the reason of non agreement of stress resultant obtained by exact solution and beam theory for a single barrel shell without edge beam of L/R=2.5 and introduced a chart for stress resultant prepared by beam arch approximation. Marshall(9) had shown it is sometimes possible to eliminate transverse moment altogether so that the stresses in the shell area reduced to those corresponding a membrane plus a constant compression. Ramaswami and Saeed (34) used Beam Method to analyse a Hyperboloid purboloda shell.

Very little work is available on the structural behaviour of the cylindrical shells with edge beams by theory for reinforced or prestressed shell.

1.4 PROPOSED INVESTIGATION:

Design of cylindrical shell by so called classical method is complicated and involves tedious computations, which, as a result, is not an attractive procedure to the engineers in practice. Beam theory, sufficiently accurate for reinforced concrete cylindrical shell, involves less complicacy. Though the method described in ASCE Manual-31 is not that complicated as other classical methods but it is not suitable for a shell with any type of combination of L/R, R/t and $\mathcal{P}_{\mathbf{K}}$. Again prestressing is desirable for optimal utilization of stressed skin structures. So there is a need for research on beam theory to explore its applicability to analylyze the stress resultant of long cylindrical shell with prestressed edge beam.

The summary of proposed investigation is:

- 1. Analysis of long cylindrical shell with edge beam
 by classical and beam theories. Then make a comparative
 study of the structural behaviour.
- 2. Parametric study of the structural behaviour of the shell by classical and beam theories.
- 3. Develop the formulations for arch analysis by strain energy method and illustrate the procedure.
- 4. Series of tables are given to design barrel shell roof by the above method.

1.5 CO-ORDINATE SYSTEMS AND SIGN CONVENTIONS:

The co-ordinate system adopted in the investigation is shown in the Figure 1.1(b). The following sign conventions are used unless otherwise mentioned (ref. fig. 2.2).

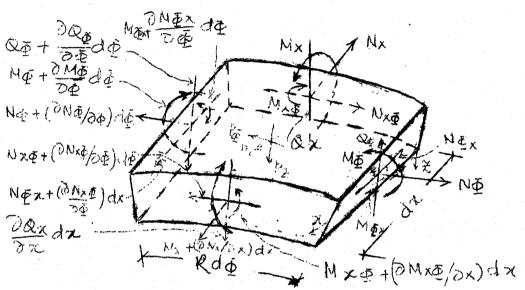


Fig. 1.2 STRESS RESULTANTS

 N_x = Longitudinal force considered positive if it is tensile.

Nø = Transverse implaneforce in **outward** direction is positive.

 $N_{x\emptyset}$ = Tangential shearing force is positive when is creates tension in the increasing direction of \emptyset and x.

Mø = Transverse moment is positive it induces tension at the bottom fiber.

u_x = Longitudinal deformation is positive if it increases
along positive x-direction.

- u_y = It is positive if it increases along positive
 y-direction.
- ur = Radial displacement is positive in it acts outwords.
- ww = Vertical deformation in downward direction is positive.
- wh = Horizontal deformation in outward direction is positive.
- W_X = Longitudinal deformation is positive when it is extensional in nature.
- w_m = Rotational deformation is positive when it has anticlock
 wise sense.

For other purpose general sign convention will be followed. Throughout the study sign convention, as mentioned above, will be followed unless other-wise mentioned.

1.6 CLASSICAL SHELL THEORY(8):

Classical shell theory developed, has the following steps or system,

- i) Primary system: Shell resisting surface load solely by direct implane stress resultant i.e. membranec state.
- ii) Errors: Forces generated by membrane theory at edges which does not satisfy the boundary conditions.
- iii) Corrections:- Corrective line loads are applied at free edges to satisfy actual boundary conditions.
- iv) Bending Solution:- These corrective line loads will cause bending of the shell for which bending solution is

required. Membrane solution is obtained considering, that, the shell is loaded at the edge by corrective line loads.

Complete solution yields by superimposing membranes solution, obtained considering shell is loaded by external loads, on bending solution, considering edge disturbance as the loading on the shell.

a) Assumptions:

- 1. Materials are homogeneous, isotropic and obey Mooke's law.
- 2. Deformation of shell produced due to loads are small.
- 3. Thickness of shell is small when compared with is is other dimensions and radii of curvature.
 - 4. Kirchoff's hypothesis holds good i.e.,
- i) Rectilinear elements normal to the middle surface of the shell remain rectilinear and normal even after deformation and do not change their lengths.
- ii) Normal stresses acting on the planes parallel to the middle surface are neglected in compresson with other stresses.
 - 5. Poisson's ratio of concrete is neglected.
- 6. Shell is assumed to be simply supported spraining between two longitudinal transcense.

b) Membrane Theory:

Membrane theory assumes external loading and the surface of this shell are continuous.

The equilibrium equations for membrane state of stress (ref. Fig. 1.2) are as follows:

$$\frac{\partial N_{x}}{\partial x} + \frac{1}{R} \frac{\partial N_{x} \emptyset}{\partial \emptyset} + q_{x}^{R} = 0$$
 (1.1)

$$\frac{1}{R} \frac{\partial N_{\emptyset}}{\partial \emptyset} + \frac{\partial N_{X}\emptyset}{\partial x} + q_{y}R = 0 \qquad (1.2)$$

$$N_{g} + q_{z}^{R} = 0 \qquad (1.3)$$

Membrane solution is obtained by integrating above solution which are given below,

i) For loading acting on the chord area of the shell $(\textbf{q}_{\mathbf{u}})$

$$N_{x} = -q_{u} R \left(\frac{1}{R}\right)^{2} \left(\cos^{2} \beta \cdot - \sin^{2} \beta \cdot\right) \frac{3}{n^{2}\pi^{2}} \sin \frac{n\pi x}{1}$$
(1.4)

$$N_{x\beta} = -q_u R \left(\frac{1}{R}\right) \cos \beta \cdot \sin \beta \cdot \frac{3}{n\pi} S \cos \frac{n\pi x}{1}$$
 (1.5)

$$N_{p} = -q_{u}R \quad \cos^{2} p \cdot \frac{1}{n} \sin \frac{n\pi x}{1}$$
 (1.6)

ii) For the loading action on the surface of the shell (q_g)

$$N_{x} = -q_{g}R \left(\frac{1}{R}\right)^{2} \cos \phi \cdot \frac{2}{n^{2}\pi^{2}} \sin \frac{n\pi x}{1}$$
 (1.7)

$$N_{x\emptyset} = -q_g R \quad (\frac{1}{R}) \quad \sin \emptyset \cdot \frac{2}{n\pi} \quad \cos \frac{n\pi x}{1}$$
 (1.8)

$$N_{\emptyset} = -q_{g}R \quad Cos \quad \emptyset \stackrel{1}{\stackrel{\bullet}{n}}Sin \quad \frac{n\pi x}{1}$$
 (1.9)

The expression of membrane displacements are given below,

i) For qu,

$$u_{x} = q_{u} R \frac{R}{Et} (\frac{1}{R})^{3} (\cos^{2} \emptyset \cdot - \sin^{2} \emptyset \cdot) \frac{3}{n^{3}\pi^{3}} \cos \frac{n\pi x}{1}$$
(1.10)

$$u_y = -q_u R \frac{R}{Et} \left(\frac{1}{R}\right)^4 \cos \theta \cdot \sin \theta \left[2 + \left(\frac{n\pi R}{I}\right)^2\right] \sin \frac{n\pi x}{I}$$
(1.11)

$$u_{r} = -q_{u} R \frac{R}{Et} \left(\frac{1}{R}\right)^{4} \left\{ \left(\cos^{2} \emptyset \cdot - \sin^{2} \theta \cdot \right) \left[2 + \left(\frac{n\pi R}{I}\right)^{2} \right] + \frac{1}{6} \left(\frac{n\pi R}{I}\right)^{4} \cos^{2} \theta \cdot \right\} \frac{6}{n^{4}\pi^{4}} \sin \frac{n\pi x}{I} (1.12)$$

ii) For
$$q_g$$
,
$$u = q_g R \frac{R}{Et} \left(\frac{1}{R}\right)^3 \cos \beta \cdot \frac{2}{n^3 \pi^3} \sin \frac{n\pi x}{1}$$
(1.13)

$$u_x = -q_g R \frac{R}{Et} \left(\frac{1}{R}\right)^4 (1 + 2n^2 \pi^2 \left(\frac{R}{1}\right)^2) \sin \emptyset^* \frac{2}{n^4 \pi^4} \sin \frac{n\pi x}{1}$$
(1.14)

$$u_{r} = -q_{g}R \frac{R}{Et} \left(\frac{1}{R}\right)^{4} \left(1 + 2n^{2} \pi \left(\frac{1}{1}\right)^{2} + \frac{n^{4} \pi^{4}}{2} \left(\frac{R}{1}\right)^{4}\right)$$

$$\times \cos \emptyset \cdot \frac{2}{n^{4} \pi^{4}} \sin \frac{n\pi x}{1} \qquad (1.15)$$

where , \emptyset = $\emptyset_{\mathbf{x}}$ - \emptyset

of loading For the expression 1. 4 to 1.15 the sinusoidal variation, along the longitudinal direction is assumed.

c) Bending Solution:

The equilibrium equations for an element on the shell (Fig. 1.2), are as follows,

$$\frac{\partial N_{x}}{\partial x} R + \frac{\partial N_{x} \emptyset}{\partial \emptyset} + Q_{x} R = 0$$
 (1.16)

$$\frac{\partial N_{\emptyset}}{\partial \emptyset} + \frac{\partial N_{\times \emptyset}}{x} R - Q_{\emptyset} + Q_{\emptyset} R = 0$$
 (1.17)

$$\frac{\partial Q_{X}}{\partial x}R + \frac{\partial Q_{\emptyset}}{\partial \emptyset} + N\emptyset + Q_{X}R = 0$$
 (1.18)

$$- \frac{\partial M \not o}{\partial \not o} + \frac{\partial M \times \not o}{\partial \times} R + Q \not o R = 0$$
 (1.19)

$$- \frac{\partial M_{x}}{\partial x} R - \frac{\partial M_{x}}{\partial x} + Q_{x}R = 0$$
 (1.20)

In the above sets of equations total number of equations is 6 and number of unknowns is 9. To solve this statically indeterminate systems the other conditions are taken from stress-displacement relationship which are given below;

$$N_{x} = Et \frac{\partial u_{x}}{\partial x}$$
 (1.21)

$$N_{\mathbf{g}} = \text{Et} \left(\frac{1}{R} \frac{\partial u_{\mathbf{y}}}{\partial \mathbf{g}} - \frac{u_{\mathbf{r}}}{R} \right) \tag{1.22}$$

$$N_{x\not 0} = \frac{Et}{2} \left(\frac{\partial u_y}{\partial x} + \frac{1}{r} \frac{\partial u_x}{\partial \emptyset} \right) \qquad (1.23)$$

$$M_{x} = -\frac{Et^{3}}{12} \frac{\partial^{2} u_{r}}{\partial x^{2}}$$
 (1.24)

$$\mathbb{M}_{\emptyset} = -\frac{\mathbb{E}t^{3}}{12} \left(\frac{1}{\mathbb{R}^{2}} \frac{\partial u_{y}}{\partial \emptyset} + \frac{1}{\mathbb{R}^{2}} \frac{\partial^{2} u_{r}}{\partial \emptyset^{2}} \right) \qquad (1.25)$$

$$M_{x\emptyset} = \frac{Et^3}{12} \left(\frac{1}{2R} \frac{\partial u_y}{\partial x} + \frac{1}{r} \frac{\partial^2 u_r}{\partial x \partial \emptyset} \right) \qquad (1.26)$$

with the help of the equations 1.16 to 1.26 the final 8th order differential equation with respect to $\mathbf{w_r}$ obtained as,

$$\nabla^{8} u_{r} + \frac{12}{t^{2}} \frac{\partial^{4} u_{r}}{\partial x^{4}} = \frac{12}{Et^{3}} R^{3} \nabla^{4} q_{z} - \frac{2\partial^{3} q_{y}}{\partial x^{2} \partial y} - \frac{1}{R^{3}} \frac{\partial^{3} q_{y}}{\partial y}$$

$$+ \frac{1}{R} \frac{\partial^{3} q_{x}}{\partial x \partial y^{2}} \qquad (1.27)$$

Where
$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{1}{R^2} \frac{\partial}{\partial g^2}$$
 (1.27)

Homogeneous solution of eqn. (1.27) has the form

$$u_r = \sum_{n=1.3}^{\infty} A_m e^{n\pi}$$
 Sin kx, where $k = \frac{n\pi}{1}$ (1.28)

putting Eq. (1.28) on the homogeneous part of 1.27 the characteristics eqn. stands as follows;

$$R^2 \left(- \kappa^2 + \frac{m^2}{R^2} \right)^4 + \frac{12}{\tau^2} \kappa^4 = 0$$
 (1.29)

Roots of the Eq. (1.29)

$$m_{14} = \pm (\alpha \pm i\beta)$$
 (1.30)

$$m_{5,8} = \pm (\alpha' + i\beta)$$
 (1.31)

where
$$\alpha = \eta \sqrt{\frac{\sqrt{(1+ \sqrt{1})^2 + 1} + (1+ \sqrt{1})}{2}}$$
 (1.32)

$$\beta = \eta \sqrt{\frac{\sqrt{(1+\sqrt{1})^2+1}-(1+\sqrt{1})}{2}}$$
 (1.33)

$$\alpha' = \eta \sqrt{\frac{\sqrt{(1-\sqrt{1})^2+1} - (1-\sqrt{1})}{2}}$$
 (1.34)

$$\beta' = \eta \sqrt{\frac{\sqrt{(1-\sqrt{)}^2+1}+(1-\sqrt{)}}{2}}$$
 (1.35)

$$V = \left(\frac{k_R}{n}\right)^2 \tag{1.36}$$

$$\eta^{8} = 3 (kR)^{4} (\frac{R}{t})^{2}$$
 (1.36)

Therefore,

$$u_{r} = \begin{bmatrix} A_{1}e^{(\alpha_{1}+i\beta_{1})} & \emptyset & + A_{2}e^{(\alpha_{1}-i\beta_{1})} & \emptyset \\ & = + A_{3}e^{(\alpha_{1}^{i}+i\beta_{1}^{i})} & + A_{4}e^{(\alpha_{1}^{i}-i\beta_{1}^{i})} & \emptyset \\ & + A_{5}e^{(-\alpha_{1}^{i}+i\beta_{1}^{i})} & + A_{6}e^{(-\alpha_{1}-i\beta_{1}^{i})} & (1.37) \\ & + A_{7}e^{(-\alpha_{1}^{i}+i\beta_{1}^{i})} & + A_{8}e^{(-\alpha_{1}^{i}-i\beta_{1}^{i})} & \end{bmatrix} \text{ Sin } kx$$

Expressing A_1 A_8 in the form $\mathbf{a} + \mathbf{b}$ and measuring the angle \emptyset ' from the crown the stress resultant and the displacement can be expressed by following expression,

$$F = \lambda (a B_1 - b B_2) \Psi_1 - (a B_2 + b B_1) \Psi_2 + (c B_3 - d B_4) \Psi_3$$
$$- (c B_4 + d B_3) \Psi_4 \qquad (1.38)$$

where.

$$\Psi_1 = \cos \beta \emptyset \cdot \cosh \alpha \emptyset$$

$$\Psi_2 = \sin \beta \, \emptyset' \, \sinh \alpha \, \emptyset'$$

$$U_{3} = \cos \beta' \not \circ ' \cosh \alpha' \not \circ ' \qquad (1.39)$$

$$\Psi_4 = \sin \beta' \sinh \alpha' \phi'$$

For different stress resultant and displacements, the shell coefficients are given in Table A.1.1 and A.1.2 in Appendix-1. Particular solution: The particular integral of Eq. 1.27 may be taken as

$$q_x = q \cos (\phi_k - \phi) \sin kx$$
 (1.40)

Substituting of the Eq. 1.40 in Eq. 1.27 gives

$$C = \frac{1}{Et} \frac{R^2 k^4 + 4 k^2 + 2/R^2}{(R^2 t^2/12) (k^2 + 1/R^2)^{4} + k^4}$$
 (1.41)

Where
$$q = \frac{4}{\pi} q_g$$

If the loading is qu then the particular solution is

$$q_x = q$$
 Sin kx, where $q = \frac{4}{\pi} q_u$

1.7 BEAM THEORY:

Analysis of long cylindrical shell by beam theory has essentially two steps, i) beam analysis, and ii) arch analysis.

a) Assumptions:

Beam theory is based on the same assumptions as classical theory given in article 1.6(a), except that certain factors are assumed to be negligible and they are:

- i) Relative displacements within each transverse cross sections.
- ii) Longitudinal bending moments ${\tt M}_{_{\rm X}}$ and radial shear forces ${\tt Q}_{_{\rm X}}$ on the shell.
 - iii) Torsional moments ($M_{\chi \emptyset}$)
 - iv) Strain from inplane shearing force.

b) Beam Analysis:

The object of beam analysis is to calculate $N_{\mathbf{X}}$, $N_{\mathbf{X}}$ Ø.

Fig. 1.3: Curved Beam Section

Moment of imertia of the section is;

$$I_c = k_I R^{3}t$$
, $k_I = \emptyset_k + \sin \emptyset_k$ (Cos $\emptyset_k - 2 \sin \emptyset_k / \emptyset_k$)
$$(1.44)$$

First moment of area at any section Ø is;

$$Q = k_Q R^2 t$$
, $k_Q = (\sin \beta - \frac{\beta}{\beta_k} \sin \beta_k)$ (1.45)

Distance of the top fiber from C.G. = Z_{yt} = Z_c

Distance of the bottom fiber from C.G. = Z_{yb} = h - Z_c

Therefore

$$N_{x} \text{ at top} = q_{0} 1^{2} \frac{t Z_{yb}}{I_{c}} \left[\frac{1}{2} \frac{x}{1} \left(1 - \frac{x}{1} \right) \right]$$
 (1.46)
$$N_{x} \text{ at bottom} = q_{0} 1^{2} \frac{t Z_{yb}}{I_{c}} \left[\frac{1}{2} \frac{x}{1} \left(1 - \frac{x}{1} \right) \right]$$
 (1.47)
$$(\text{tension})$$

$$N_{x}$$
 at any section = $q_0 \cdot 1 + \frac{Q}{2I} \cdot \left(\frac{1}{2} - \frac{x}{1}\right)$ (1.48)

c) Arch Analysis:

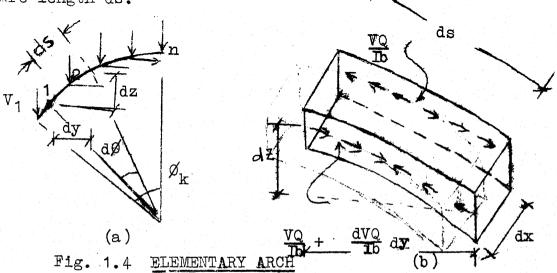
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Object of the arch analysis is to calculate Ng , Mg and Qg.

An elementary arch strip, which is under the action of q_u , q_g and specific shear, is considered. Eq.(1.45)express the magnitude of $N_{\chi 0}$ at a section \emptyset along χ .

Therefore,
$$\frac{\partial N_{x} \neq 0}{\partial x} = -\frac{q_{0,Q}}{2 I}$$
 (1.49)

So the specific shear at any longitudinal section is constant and ints vertical component will act upwards and horizontal component inwards. The arch is divided into a no. of nodes of arc length ds.



The arc length between two adjacent nodes i and i+1 is = ds (= RdØ.)

Vertical projection of $ds = dZ_i = R$ ($Cos \beta_{i+1} - Cos \beta_i$)
(1.50)

Horizontal projection of ds= $dy_i = R(Sin p_i - Sin p_{i+1})$

(1.51)

Therefore, $F_{vi}^{r} = q_u dy_i + q_g ds_i + \frac{\partial N_{x0}}{\partial x} dz_i$ (1.52)

$$F_{hi} = \frac{\partial N_x \emptyset}{\partial x} dy_i \qquad (1.53)$$

Therefore,

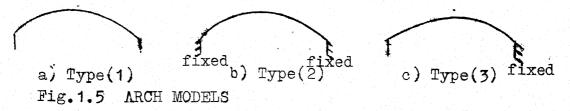
$$M_{\emptyset i} = \sum_{j=1,2}^{i} \left[F_{vj} (Y_i - Y_j) + F_{hj}(Z_i - Z_j) \right]$$
 (1.54)

$$\mathbb{N}_{\phi_{\mathbf{i}}} = \left(\begin{array}{c} \mathbf{i} \\ \mathbf{\Sigma} \\ \mathbf{j=1,2} \end{array} \right) \operatorname{Sin} \phi_{\mathbf{i}} - \left(\begin{array}{c} \mathbf{i} \\ \mathbf{\Sigma} \\ \mathbf{j=1,2} \end{array} \right) \operatorname{Cos} \phi_{\mathbf{i}}^{\mathbf{i}}$$
(1.55)

$$Q_{\emptyset_{i}} = \begin{pmatrix} i & F_{\forall j} \end{pmatrix} \cos \emptyset_{i} - \begin{pmatrix} i & F_{hj} \end{pmatrix} \sin \emptyset_{i} \quad (1.56)$$
1.8 DISCUSSION;

The boundary conditions for arch calculations is determined from the physical system of the shell itself, which can be divided mainly in 3 types which are given below:

1) Single Barrel:- The arch may be considered as axisymmetric statically determinate system, loaded by external loads and supported by specific shear, having two ends free.



- 2) Interior Shell:- The arch may be considered as axisymmetric statically inderminate system, loaded, as, type-1, with both ends fixed.
- 5) End Shell: The arch may be considered as axi-asymetric statically determinate system, loaded as type 1, with one end fixed and another end free.

In the next chapter both the methods are illustrated by numerical examples.

CHAPTER-II

ANALYSIS OF CYLINDRICAL SHELL BY CLASSICAL AND BEAM THEORIES

2.1 INTRODUCTION:

The stress resultants and the displacements of a shell can be expressed through four constants a , b , c , d which can be determined from boundary conditions. For the shell without edge beam the boundary conditions are that $N_{\rm x}$, $N_{\rm x0}$, $N_{\rm y}$ are zero at the free edge. If the shell is provided with edge beam, the boundary conditions are the edge deformations of shell are same as those of edge beam at the interface of the shell and beam.

2.2 THEORY OF EDGE BEAM:

For the shell with edge beam the compatibility of deformations are:

- 1. Compatibility of vertical deformation
- 2. Compatibility of horizontal deformation
- 3. Rotational compatibility and
- 4. Longitudinal deformation compatibility.

If the stress resultants and the deformations are denoted by the subscript 'k' then,

Vertical deformation, $w_{vs} = w_{rk} \sin \beta_k$ Horizontal deformation $w_{hs} = w_{rk} \cos \beta_k + u_{yk}$ (2.1) Rotational deformation $w_{ms} = \theta_k$ Longitudinal deformation $w_{xs} = u_{xk}$ Loading of the edge beam at the interface is the edge stress resultants of the shell edge in same magnitude but of opposite sign along with the external loads, if any, and self weight of the edge beam. If V, H, M and S are the vertical, horizontal, rotational and shear loads respectively at the edge of the shell, the edge deformations are given by:

$$w_{vb} = w_{vb} (V , H , M , S , q_b)$$
 $w_{hb} = w_{hb} (V , H , M , S , q_b)$
 $w_{mb} = w_{mb} (V , H , M , S , q_b)$
 $w_{xb} = w_{xb} (V , H , M , S , q_b)$

(2.2)

Where qb represents the external loading on the edge beam.

Again V , H , M , S , w_{vs} , w_{hs} , w_{ms} and w_{xs} are associated with a , b , c , d discussed in previous chapter and they can be expressed as functions of a , b , c , d if the loading and geometrical properties of shell and edge beam are known. Hence eqs. of 2.1 and 2.2 yield the interface compatibility as follows,

$$w_{vs}$$
 (a, b, c, d, q_s) = w_{vb} (a, b, c, d, q_b)
 w_{hs} (a, b, c, d, q_s) = w_{hb} (a, b, c, d, q_b) (2.3)
 w_{ms} (a, b, c, d, q_s) = w_{mb} (a, b, c, d, q_b)
 w_{xs} (a, b, c, d, q_s) = w_{xb} (a, b, c, d, q_b)

where q_b and q_s are the total loading system on shell and edge beam. The four equations of 2.3 will now determine the four constants.

2.3 DEFORMATIONS OF EDGE BEAM:

Details of loading on edge beam are shown in Fig. 2.1. The expressions of deformations are given below;

a) Vertical Deformation:

$$w_{vV} = \frac{1}{m^{2}} \left[\frac{1}{m^{2} E_{b} I_{y}} + \frac{(\pi - 2) e_{h}^{2}}{2 G_{b} J} \right] V = e_{b11} V$$

$$w_{vS} = -\frac{1}{m^{3}} \frac{1}{E_{b} I_{y}} e_{v} S = c_{b12} S$$

$$w_{vH} = -\frac{1}{m^{2}} \frac{\pi - 2}{2 G_{b} J} e_{h} e_{v} H = c_{b13} H$$

$$w_{vM} = -\frac{1}{m^{2}} \frac{\pi - 2}{2 G_{b} J} e_{h} M = c_{b14} M$$

$$(2.5)$$

$$w^{\Delta p} = w^{\Delta A} + w^{\Delta H} + w^{\Delta M} + w^{\Delta}$$

Where the first subscript indicates the direction of deflection and the second indicates the force causing it.

b) Longitudinal Deformation:

$$w_{xV} = -\frac{1}{m^2} \frac{1}{E_b I_y} e_v V = c_{b21} V$$

$$w_{xS} = -\frac{1}{m} \left[\frac{1}{A_b} + \frac{e_v^2}{E_b I_y} + \frac{e_h^2}{E_b I_z} \right] S = c_{b22} S$$

$$w_{xH} = -\frac{1}{m^3} = \frac{1}{E_b I_z} e_{hH} = c_{b23}H$$
 (2.6)

$$w_{xM} = 0$$

$$w_{xb} = w_{xy} + w_{xs} + w_{xH}$$

e) Horizontal Deformation:

$$w_{h \cdot V} = -\frac{1}{m^2} \frac{\pi - 2}{2G_b J} e_h e_v V = c_{b31} V$$

$$w_{hS} = -\frac{1}{m^3} \frac{1}{E_{b}I_{z}} e_{h} S = c_{b32}S$$

$$w_{hH} = \frac{1}{m^2} \left[\frac{1}{m^2} \frac{1}{E_b I_z} + \frac{\pi - 2}{2G_b J} e_v e_h \right] H = c_{b33} H (2.7)$$

$$w_{hM} = -\frac{1}{m^2} \frac{\pi - 2}{2G_h J} e_v M = c_{b34} M$$

$$w_{hb} = w_{hV} + w_{hS} + w_{hH} + w_{hM}$$

d) Rotational Deformation:

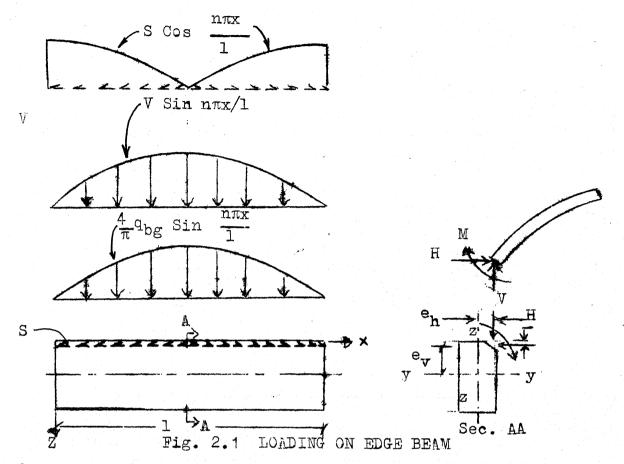
$$\mathbf{w}_{\mathbf{m}\cdot\mathbf{V}} = - \frac{1}{\mathbf{m}^2} \frac{1}{\mathbf{E}_{\mathbf{b}}\mathbf{I}_{\mathbf{y}}} \mathbf{e}_{\mathbf{v}} \mathbf{V} = \mathbf{e}_{\mathbf{b}41} \mathbf{V}$$

$$w_{m S} = 0$$

$$w_{mH} = -\frac{1}{m^2} \cdot \frac{\pi - 2}{2G_{h} J} e_{v} H = e_{b43}^{H}$$
 (2.8)

$$w_{m H} = \frac{1}{m^2} \frac{\pi - 2}{2 G_{b} J} M = c_{b44} M$$

$$\mathbf{w}_{\mathrm{mb}} = \mathbf{w}_{\mathrm{mV}} + \mathbf{w}_{\mathrm{mH}} + \mathbf{w}_{\mathrm{mM}}$$



2.4 DEFORMATION OF EDGE BEAM DUE TO PRESTRESSING FORCE:

Prestressed cable in a vertical plane will cause deformation in vertical plane only. These deformations are vertical and honightudial deformations. The edge beams are prestressed by curved cable (tendons) having parabolic profile. The effect of prestressing the edge beams is equivalent to the particular cases of loading which are described below and shown in Fig. 2.2.

1. Continuous distributed upward pressure(wbe) due to the change of curvature of the cable profile and if magnitude is given by,

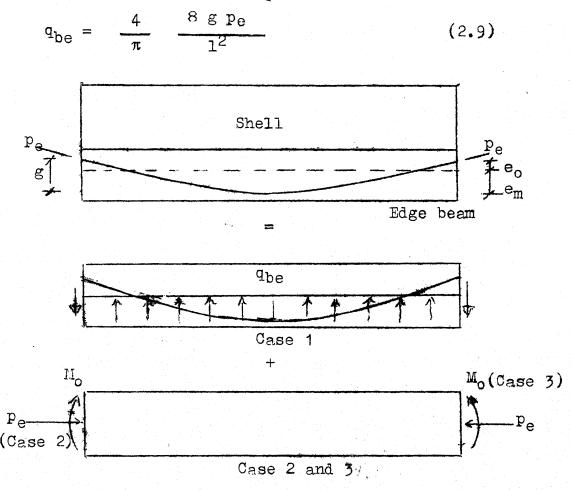


Fig. 2.2 EFFECT OF PRESTRESSING FORCE

2. End moment due to eccentricity of anchorage is given by,

$$M_{Q} = \frac{4}{\pi} p_{e} q_{O} \qquad (2.10)$$

3. Axial thrust applied at both ends of the centre of gravity of the beam section given by, axial thrust = $\frac{4}{\pi} \frac{p_e}{A_b}$ (2.11)

Vertical deformation caused by prestressing force,

$$w_{vp} = -\frac{4}{\pi} \left(\frac{1}{m^4} \frac{8g}{E_b I_v} - \frac{t^2}{m^2} \frac{e_o}{E_b I_v} \right) \frac{p_e}{1^2}$$
 (2.12)

and longitudinal deformation,

$$w_{xp} = \frac{4}{\pi} \left(\frac{1}{A_b} + \frac{e_0 e_v}{E_b I_y} - \frac{1}{m^2} \frac{8g e_v}{1^2} \frac{1}{E_b I_y} \right) p_e$$
 (2.13)

The coefficient $4/\pi$ is used in all the Eqs. 2.12 and 2.13 as a Fourier coefficient of uniformly distributed load converted into sine series.

2.5 ANALYSIS OF SHELL BY CLASSICAL THEORY:

The theory described in article 1.6 is the theoretical aspect of cylindrical shell analysis by the method described in ASCE Manual-31. A series of comprehensive tables are given in ASCE Manual-31 for barrel roof analysis. In this investigation shells are analysed using those tables taking n=1.

All the coefficients given in the table of ASCE Manual-31 for the loads assuming sinusoidal variation of loads along the longitudinal direction. The method to analyse the shell is given below,

Membrane Analysis:-
$$N_{x} = \frac{4}{\pi} \left(\frac{1}{R}\right)^{2} q R^{c}_{1m} Sin mx$$

$$N_{x} = \frac{4}{\pi} \left(\frac{1}{R}\right)^{2} q R c_{2m} Cas mx$$
(2.14)

$$N_{\emptyset} = \frac{4}{\pi} q R c_{3m} Sin mx$$

Where q is the loading For q = q_g , c_{1m} , c_{2m} , c_{3m} are given in column 7, 8, 9 of Table 1 B and for q = q_g column 1, 2, 3 gives the above coefficients. If $c_s = 1^4/(R^3 t E_s)$, then the membrane displacement can be obtained from the following expression,

i) For
$$q_g$$
 $w_{VM} = q_g R c_s \left[\left(\frac{2R}{\pi l} \right)^2 + \frac{2}{\pi^4} + \left(\frac{R}{l} \right)^4 c_{w1m} \right] Sin mx$
 $w_{hm} = q_g R c_s \left[\left(\frac{R}{l} \right)^4 c_{w2m} Sin m x \right]$

(2.15)

 c_{w1m} and c_{w2m} are given in column 10 and 11 of Table 1 B.

ii) For
$$q = q_u$$

$$w_{vm} = q_{u}R c_{s} \left[\left\{ 1 + \frac{1}{2} (\pi R/1)^{2} + (\pi R/1)^{4}/12 \right\} c_{w1m} \right] \sin mx$$

$$w_{hm} = q_{u}R c_{s} \left[(R/1)^{4} c_{w2m1} + \left\{ 1 + (\pi R/1)^{2}/2 + (\pi R/1)^{4}/12 \right\} \chi$$

$$c_{w2m2} \int \sin mx \qquad (2.16)$$

 $c_{\rm w1m},~c_{\rm w2m1}~$ and $c_{\rm w2m2}$ are given in column 4 , 5 and 6 of Table 1 B and the longitudinal deformation,

$$w_{xm} = \frac{1}{E_s t} (Nx) / (1/R)^2 = \frac{1}{E_s t} f_{osm}$$
 (2.17)

Errors and Correction: To satisfy the boundary condition i.e. $N_{\rm X}$, $N_{\rm X}$ 0 , N0 at shell edge are zero, a set of line loads

should be applied at the edge which is equal in magnitude but opposite in sign of N_{xk} , $N_x p_k$ and $N p_k$. Correction forces are resolved into vertical and horizontal components and they are:

$$V_{L} = - N_{\not 0k}$$
 Sin $\not 0$ Sin m x

 $H_{L} = - N_{\not 0k}$ Cos $\not 0$ Sin m x

 $S_{L} = - N_{x\not 0k}$ Cos m x

Bending Analysis:-

The correction line loads will disturb the membrane state of shell and bending of shell will occur. The bending analysis of shell is done in following way.

Edge deformations due to correction line loads

$$w_{vL} = c_{s} \left[c_{v1} \quad V_{L} + c_{v.2} \quad S_{L} + c_{v3} \quad H_{L} \right] \text{Sin mx}$$

$$w_{xL} = \left(\frac{1}{R} \right)^{2} \frac{1}{tE_{s}} \left[c_{1} \quad V_{L} + c_{12} \quad S_{L} + c_{13} \quad H_{L} \right] \text{Sin mx} \qquad (2.19)$$

$$w_{ml} = \frac{R^{2}}{E_{s}I} \left[c_{m1} \quad V_{L} + c_{m2} \quad S_{L} + c_{m3} \quad H_{L} \right] \quad \text{Sin mx}$$

While c_s c_{h1} V_L + c_{h2} S_L + c_{h3} H_L Sin mx Where c_{v1} , c_{v3} , c_{v2} , c_{hi} , c_{h3} , c_{m1} , c_{m3} and c_{m2} are given in column 1,4,7;4,5,8;3,6 and 9 of Table 2B and c_{11} , c_{13} c_{12} are given in column 1,5 and 9 of Table 2A.

When the shell is provided with edge beam a set of forces, say v_b , H_b , s_b and M_b , will generate at the interface. The deformations due to the interface forces are:

where c_{m1} , c_{m3} , c_{m2} , c_{v4} , c_{h4} and c_{m4} are given in column 3, 6, 9, 10, 11 and 12 of Table 2 B and c_{14} in column 4 in Table 2 A. If positive interface load acts on the Shell

then negetive load will act on the beam, then the deformation compatibility at mid span of shell edge yields,

$$w_{vm} + w_{vL} + w_{vbl} - w_{vb} = 0$$
 $w_{xm} + w_{xL} + w_{xbl} - w_{xb} = 0$
 $w_{hm} + w_{hL} + w_{hbl} - w_{hb} = 0$
 $w_{mL} + w_{mbl} - w_{mb} = 0$

(2.21)

(3rd Subscript 1 indicates deformations due to loading)

$$\begin{bmatrix} c_{s} c_{v1} + c_{b11} & c_{s} c_{v2} + c_{b12} & c_{s} c_{v3} + c_{b13} & \frac{c_{s} c_{v4}}{R} & c_{b14} \\ (1/R)^{2} \frac{1}{E_{s} t} c_{11} + c_{b21} & (\frac{1}{R})^{2} \frac{1}{E_{s} t} c_{12} + c_{b22} & (\frac{1}{R})^{2} \frac{1}{E_{s} t} c_{13} + c_{b23} & (\frac{1}{R})^{2} \frac{c_{14}}{E_{s} t} \\ c_{s} c_{h1} + c_{b31} & c_{s} c_{h2} + c_{b32} & c_{s} c_{h3} + c_{b33} & \frac{c_{s} c_{h4}}{R} + c_{b34} \\ \frac{R^{2}}{E_{s} I} c_{m1} + c_{b41} & \frac{R^{2}}{E_{s} I} c_{m2} & \frac{R^{2}}{E_{s} I} c_{m3} + c_{b43} & \frac{c_{s} c_{m4}}{R} + c_{b44} \end{bmatrix}$$

$$(2.22i)$$

Where,

$$C_1 = w_{vm} - w_{v1} + w_{vbg} + w_{vbp}$$
 $C_2 = -w_{xm} - w_{x1} + w_{xbg} + w_{xbp}$
 $C_3 = -w_{hm} - w_{h1}$ and $C_4 = -w_{m1}$ (2.23)

wvbg = vertical deformation due to self weight of the beam

$$= c_{b11} b d V$$
 (2.24)

 w_{xbg} = longitudinal deformation due to the self weight of the beam = c_{b21} b d V (2.25)

Analysis of shell by using the formulae direct from the solution of differential equation of shell involves the evaluation of the constants a, b, c and d as basic unknowns whereas method using the tables takes the interface forces as basic knowns, which are determined from the boundary conditions. Once V_b , S_b , H_b , and M_b are known, the total correction loads;

$$V = V_{L} + S_{b}$$
, $S = S_{L} + S_{b}$, $H = H_{L} + H_{b}$, $M = M_{L} + M_{b}$ (2.26)

Stress resultants due to bending;

$$N_{x} = (e_{11} V + c_{12} S + c_{13} H + c_{14} \frac{M}{R}) \times (1/R)^{2}$$

$$N_{x} \emptyset = (c_{21} V + c_{22} S + c_{23} H + c_{24} \frac{M}{R}) \times (1/R)$$

$$N_{\emptyset} = c_{31} V + c_{32} S + c_{33} H + c_{34} M/R$$

$$M_{\emptyset} = (c_{41} V + c_{42} S + c_{43} H + c_{44} M/R)R$$

$$(2.27)$$

Total stress resultants can be obtained by super-imposing membrane stress resultants obtained from Eqs. of 2.14 to bending stress resultant obtained from Eqs. of 2.27.

It is general practice for single barrel shell to assume that edge of the shell does not offer any restraint against horizontal and rotational deformations, so Eqns. of 2.22 reduces to 2 simultaneous equations. For the interior shell of a multi barrel group w_h = 0 and w_m = 0 while for single shell H = 0; M = 0. Numerical computations of the method is illustrated by the fallowing example.

2.6 EXAMPLE 1

1 = 40 m, R = 12 m, t = 0.12 m, b = 0.15 m, d = 2.0 m. $\emptyset_k = 40^0$, $e_0 = 0.5 \text{ m}$, $e_m = 0.7 \text{ m}$, $p_e = 400,000 \text{ kg}$.

Loading:

- i) Live load on shell = 75 kg/m^2 (acting on surface of shell)
- ii) Load from lime terracing = 45 kg/m² (acting on surface of shell)
- iii) Self weight of shell = 2400 x 0.12 = 360 kg/m²

 Total load on the shell = 480 kg/m²

 Grade of concrete of shell and beam is M₃₅₀.

 The shell is Single barrel

a) Membrane Analysis:-

Membrane stresses are calculated from Eqs. of 2.14 and are given in Table 2.1.

Ø	N _x (kg/m)	N _x ø (kg/m)	Nø (kg/m)	Mø (kg/m)
40	-14040	0	-6234	0
30	-13820	-2297	-6139	0
20	-13190	-4524	-5858	0
10	-12160	- 6614	-5399	0
0	-10750	-8503	-4775	Ω

Table 2.1 MEMBRANE STRESS RESULTANTS

b) Errors and Corrections:

Errors are Ngk = - 4775 kg/m and Nxgk = - 8503 kg/m . The vertical and horizontal component shear correction line loads are calculated from Eqs. of 2.14.

 $V_L = 3070 \text{ kg/m}$

 $H_L = 3658 \text{ kg/m}$

 $S_{T_i} = 8503 \text{ kg/m}$

c) Bending Analysis:

i) Deformation of shell edge at mid span:- It is general practice to ignore the restraint offered by the edge

beam towards the horizontal and rotational deformation as the beam is slender one. So for this problem H_b and M_b may be taken as zero. From Eqs. 2.15, 2.17 and 2.19.

$$w_{vm} + w_{v1} = c_s 54500$$

$$w_{xm} + w_{x1} = (\frac{1}{R})^2 \frac{1}{E_s t} 26130$$

ii) Deformation of Edge Beam:- The contribution of e_h towards the beam deformation is very small and hence ignored. From Eqs. of 2.5

Words + w_{vbs} = c_{b1} (- 0.12319) v_b + $\frac{h}{l}$ 0.19351 Sb

From Eqs. of 2.25

$$w_{\rm vbg} = c_{\rm b1} - 112.9$$

From Eq. 2.12

$$w_{\text{vbp}} = c_{b1}$$
 1.2548 e_{m} - 0.2923 e_{0}

From Eqn of 2.6

$$w_{xvb} + w_{xbs} = c_{b2}$$
 (0.60793) $V_b - \frac{d}{L}$ 1.2732 S_b $w_{xbg} = c_{b2}$ 557.3

From Eq. 2.13

iii) Compatibility of Deformations:- Vertical deformation compatibility from Eq. 2.21,

42.38
$$V_b$$
 + 0.9561 S_b - $\frac{t}{b}$ ($\frac{R}{d}$)³(-0.1231 V_b + $\frac{d}{1}$ 0.19351 S_b + 112.9 + (1.2548 e_m - 0.2923 e_o) P_e/L^2) + 54500 = 0 (2.28)

Longitudinal compatibility from Eq. 2.21

$$-0.216 \ V_{b} - 0.1668 \ S_{b} - \frac{t}{b} \left(\frac{R}{d}\right)^{2} (0.60793 \ V_{b} - \frac{d}{1} 0.12732 \ S_{b}$$

$$-557.3 + (1.2732 \ d_{t} 1.4478 \ e_{o} - 6.1923 \ e_{m}) \frac{p_{e}}{L^{2}}) + 26130 = 0$$

$$(2.29)$$

d) Total Stress Resultants:

Eqs. 2.28 and 2.29 gives $V_b = -1226$ and $S_b = -15980$ Therefore, final correction line loads from eqs. of 2.26,

$$V = 1844$$
 , $S = -7477$ and $H = 3658$

Bending stress resultants are calculated from Eqs. of 2.27 and adding these with membrane solution total stress resultants are obtained and given in Table 2.2.

Table 2.2 TOTAL STRESS RESULTANTS

Ø.	N _X (kg/m) (at x=1/2)	N _x ø (kg/m) (at x = 0)	Ng (kg/m) (at x= 1/2)	Mø (kgr/m) (at x=1/2)
40	- 51180	O	-5393	735
30 [°]	-43180	-10100	-4919	695
20	-22950	-17090	-367 4	991
10	-27290	-19610	- 2100	865
0	-44570	-19720	-619	Q

Longitudinal stress at the bottom of the beam

$$\frac{N_{xb}}{b} = \frac{1^2}{b d^2} \left[0.60793 \left(\frac{4}{\pi} b d \sqrt{-V_b} \right) + 0.6336 \frac{d}{1} S_b + \frac{4}{\pi} \left(-d + 1.4478 e_0 - 6.1923 e_m \right) \right]$$
 (2.30)
$$= -297000 \left(\frac{kg}{m^2} \right)$$

The variation of stress resultants are shown in Fig. 2.5 (a) and (b).

2.7 ANALYSIS OF SHELL BY BEAM THEORY:

Theoretical aspect of beam theory has been described in the article. 7. The beam method may be devided into two parts as beam and arch analysis. To illustrate the numerical computations of the method the same shell of example 1 is analysed below.

a) Beam Analysis:

The transverse section of the shell along with the edge beam is assumed to constitue an uncracked section of a beam whose geometrical properties are shown in Fig. 2.3.

Chord length, $l_c=2R\sin\ \phi_k=15.4286\ m$ Depth of the section, $h=R\ (1-\cos\phi_k\)+d=4.8075\ m$ Area of the Section, $A_S=2\ (A_C+A_B\)=2.6106\ m^2$

Distance of the C.G. from crown, $Z_c = \frac{2(A_c Z_c + A_b Z_b)}{A_s} = 1.6077 \text{ m}$

Moment of inertia of the section, $I_s = 2$ ($I_b + I_c + A_b Z_{rb}^2 +$

$$\Lambda_{\rm c} Z_{\rm rc}^2$$
) = 5.3950 m⁴

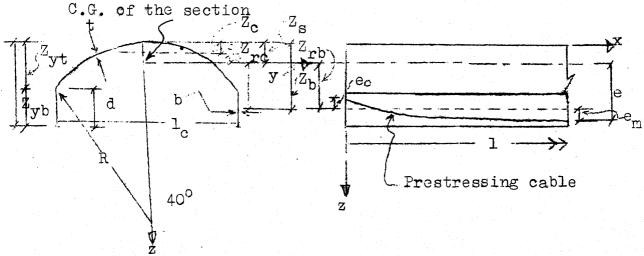


Fig. 2.3 GEOMETRICAL PROPERTIES OF ASSUMED BEAM SECTION OF EXAMPLE 1

First moment of area of any section, β .

- i) for shell portion, Q = 2Rt Z_S R (1-Sin ϕ_k/ϕ_k)
- ii) for beam portion, $Q = \int_{Z}^{Z_{yb}} 2b \ Zdz = b(Z_{yb}^2 Z^2)$

Co-ordinates of any point on the shell,

$$Z_i = R (1 - Cos \emptyset_i) - Z_s$$

$$Y_i = R \sin \phi_i$$

Total load U.D.L. on the shell on the chord area,

$$q_o = \left[(A_g \sqrt{+q_g}) / / \sin / k + q_u \right]$$
 $l_c = 8277.4946 \text{ kg/m}$

Bending moment at the mid span,

$$M = \frac{q_0 1^2}{8} = 1655.2 \times 10^3 \text{ kg. m.}$$

Shear force at the end section,

$$S_f = \frac{q_{tL}}{2}$$
 165.52 x 10³ kg.

 $N_{\rm X}/{\rm t}$ at the top of the section

$$= \frac{p_e}{\Lambda_S} - \frac{p_e e}{I_S} Z_{yt} + \frac{M}{I_S} Z_{yt} = 107700$$

 N_x/b at the bottom of the section,

$$= -\frac{p_e}{\Lambda_s} - \frac{p_e e}{I_s} Z_{yb} + \frac{M}{I_s} Z_{yb} = -702635$$

$$N_{x\emptyset} \text{ at any section } \emptyset' = \frac{S_f Q}{2 I_s}$$

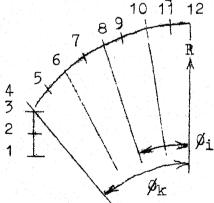


Fig. 2.4 NODES ASSUMPTIONS FOR ARCH ANALYSIS

The section is a symmetric section, so the stresses of the half section is given in the Table 2.3

Table 2.3 T_x AND $N_x \not o$ OF EX.2 (BY BEAM METHOD)

Nodal	ø٠	Co-ordi	nates	Q	$\sigma_{\rm x}$	$N_{\mathbf{x}\emptyset}$
Points		Z	Yi	((kg/m^2) x = L/2)	(kg/m) (at x=L/2)
1	_	3.2028	7.7136	0.000	107700	0000
2	-	2.2028	7.7136	0.8099		5218
3	_	1,2028	7.7136	1.320		8504
4	40	1.2028	7.7136	1.320		8504
5	35	0.5644	6.8832	1.540		9921
6	30	0.0028	6.0	1.6093		10370
7.	25	-0.4804	5.0712	1.5461		9963
8	20	-0.8816	4.1040	1.3732		8845
9	15	-1.196	3.1056	1,1089		7146
10	10	-1.4228	2.0832	0.7775		5010
11	5	-1.5596	1.0464	0.4002		2570
12	0	-1.6077	0.0000	0.0000	702635	0000

b) Arch Analysis:

An elementry arch strip of width dx is considered at midspan, which containts the shell and the edge beam together and divided into 12 modal points as shown in Figure 2. The details of arch calculations are given below in tabular form Table 2.4 contains ds, dy, dz, $\frac{\partial N_{x0}}{\partial x}$. Table 2.5

contains $\frac{\partial N_x \not o}{\partial x}$ dz , q_g ds , q_u dy , F_v , F_h (= $\frac{\partial N_x \not o}{\partial x}$ dy) and Table 2.6 contains $N_{\not o}$, $M_{\not o}$, distance of the nodes from the top. It is assumed that nodes 4 and 12 have half of the segmental length and 1 $\stackrel{?}{\sim}$ 3 have segmental length $\frac{1}{4}$ d where 2 has d/2.

Table 2.4 ARCH ANALYSIS

Nodes	$\delta \mathbb{N}_{\mathbf{x}}$	₫ \$	đу	₫z
*	gx			
1	000.000	0.5 000	0,00000	0.50000
2	260.8898	1.0000	0.0000	1.00000
3	425.7712	0.5000	0.00000	0.50000
4	424.7712	0.5236	0.88280	0.39528
5	494.9012	1.0472	0.85752	0.60048
6	517.4841	1.0472	0.90660	0.52344
7	526.5362	1.0472	0.94884	0.44244
8	441.0083	1.0472	0.98364	0.35808
9	356.7043	1.0472	1.01120	0.27096
0	250.0121	1.0472	1.03920	0.18168
1	128.9848	1.0472	1.04280	0.09132
2	000.000	0.5236	0.52334	0.01140

Table 2.5 ARCH ANALYSIS

Nodes	₩ [*] g d _s (1)	w _u dy (2)	$\frac{\partial N_{x} \emptyset}{\partial x} dz$	(1)+(2)+(3)	(dn dy)
1	213.6288	0	000.000	213.6288	000.0000
2	427.2576	0	- 11.7789	415.4787	000.0000
3	427.2576	0	- 45.4222	381.8325	000.0000
4	427.2576	0	- 96.6527	330.6049	207.0933
5	427.2576	0	-157.9163	269.3413	424.3885
6	427.2576	0	-232.9607	194.2969	469.1511
7	427.2576	0	-270.8705	156.3871	499.5986
8	427.2576	0	-297.1789	130.0787	434.7934
9	213.6288	0	-167.64932	74.7334	360.6997
10	180.0000	0	-212.8856	-32.8856	257.7425
11	360.0000	o	-260.8898	99.1102	134.5054
12	180.0000	0	000.0000		000,000

 w_g for shell 408 kg/m² and for beam 360 kg/m² $w_q = q_q$ $w_u = q_u$

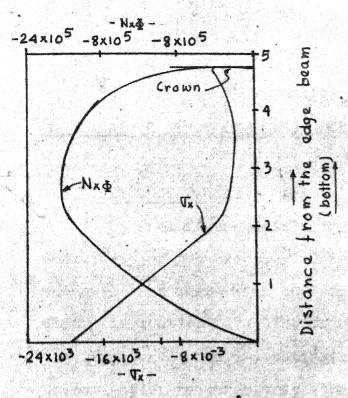
Table 2.6 Nø AND Mø (BY BEAM METHOD)

Nodes	hį		Mø kg m/m (at x= L/2)
1	4.8110	0000	0000.0000
2	3.8110	0000	0000,0000
3	2.8110	0000	0000,0000
4	2.8110	-1430	0000,0000
5	2.1726	-1580	1635.1483
6	1.5954	-1804	2961.8528
7	1.1248	-2136	4094.0868
8	0.7236	-2387	4997.2154
9	0.3560	-2591	5683.6779
10	0.1754	-2712	6160.6666
11	0.0456	-2794	6428.3211
12	0.0000	-2887	6508.4988

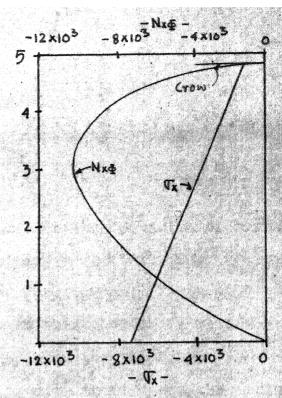
2.8 DISCUSSION:

It is evident from the articles 2.6 and 2.7 that there is no agreement between the stress resultants evaluated by beam method and classical method. A comparative study for the structural behaviour determined by both the theories is necessary to get the idea about the error involvement in beam theory. The influence of the different shell parameters,

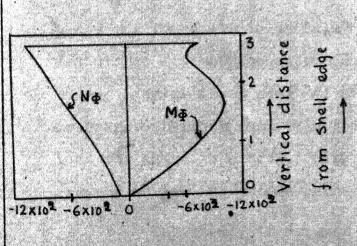
prestressing force and eccentricities of the cable should be investigated. In next chapter(3) a comparative study of the structural behaviour of the shell of example 2 and parametric study are done.



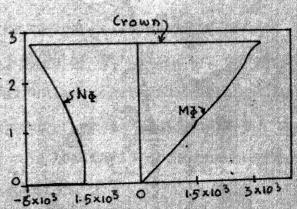
a) Tx & Nx& by classical theory



c) Tx & Nx4 by beam theory



b) No & Mo by classical theory



d) No L Mo by beam theory

Fig. 2.5 VARIATION OF STRESS RESULTANTS (Ex.-1)

CHAPTER-III

COMPARATIVE STUDY OF STRUCTURAL BEHAVIOUR BY BEAM AND CLASSICAL THEORIES

3.1 INTRODUCTION:

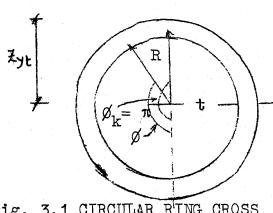
From figure 225it is found that the variation of stress resultants of the example-1 determined in section 2.6 by classical theory do not agree with the stress resultants determined by beam theory in section 2.7. The object of this chapter is to investigate our the structural behavi-/ estimated by both the theories.

Investigation by parametric study is done to find out the influence of different shell parameters and the prestressing force on the structrual behavior.

3.2 ANALYSIS OF SHELL IN FULLY MEMBRANE STATE BY BEAM THEORY (2)

The shell is in fully membrane state when it offers resistance against the external forces by it's implane stresses only. A cylindrical shell with circular generatrix attains fully membrane condition when it's transverse cross section is continued i.e. the shape is circular ring and loading is continuous with respect to its transverse cross section.

The shell, to analyse by beam theory, is assumed as a beam spanning intthe direction of it's axis of rotation with a hollow circular transverse cross section.



Loading on the shell = $2\pi Rq_g$ Sin mx Moment of Inertia, $I_e = \pi R^3 t$ First moment of area at any section

 $= 2R^2 t \sin \phi$

Bending moment at any lognitudiral secti section, $M = 2 \pi R q_g \cdot \frac{1}{m^2}$ Sin mx

Fig. 3.1 CIRCULAR RING CROSS Section, SECTION Shear force $S_f = -\frac{1}{m} 2 \pi R q_g \cos m\bar{x}$

 $Z_{yt} = R \cos (\pi/2 - \emptyset) = -R \cos \emptyset$

Longitudinal stress, $T_x = -2 \times Rq_g = \frac{1}{m^2} \sin mx = \frac{-R \cos \emptyset}{\pi R^3 t}$ $Z = -2^{\pi} Rq_g \frac{1}{m} \cos mx \frac{2 R^2 t \sin \phi}{2 \pi t R^3 t}$ (3.1) Shear stress,

Trans verse force, $N_{\emptyset} = -q_g \sin mx$ (- R Cos \emptyset_k) Substituting $N_x = \sqrt{x}t$, and $N_x = 7$ t and $m_x = n\pi/1$ eqn. of 3.1.

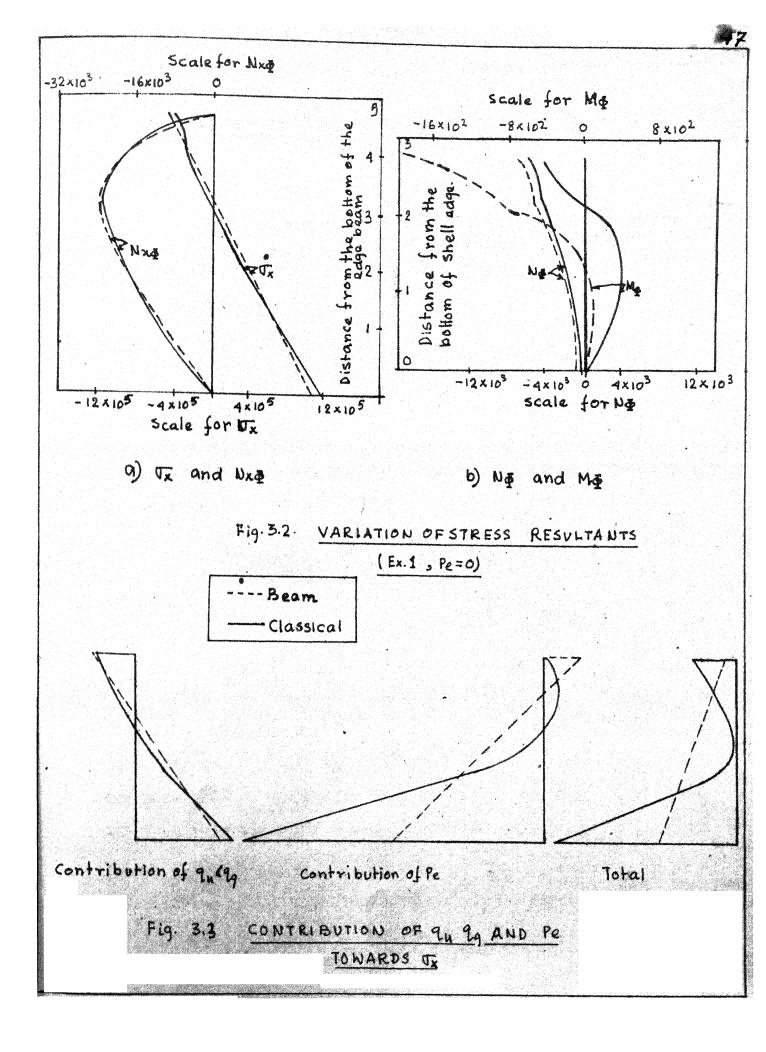
$$N_{x} = q_{g} R \left(\frac{1}{R}\right)^{2} \cos \emptyset. \frac{2}{n^{2} \pi^{2}} \sin \frac{n\pi x}{1}$$

$$N_{x}\emptyset = -q_{g} R \left(\frac{1}{R}\right) \sin \emptyset \frac{2}{n\pi} \cos \frac{n\pi x}{1}$$

$$N_{\emptyset} = q_{g} R \cos \emptyset \sin \frac{n\pi x}{1}$$

$$(3.2)$$

If $\emptyset_k = \pi$ is substituted then eqns. 1.7, 1.8, 1.9 would give exact same expressions like eqa. of 3.2



3.3 COMPARATIVE STUDY OF THE STRUCTURAL BEHAVIOUR

our

By the term structural behavi-/it is meant the response of the shell by stress resultants and the deformations. As the investigation mainly aims to explore the feasibility of the aplicability of beam theory to design cylindrical shell the following parameters are considered for structural behaviour:

 T_x , N_0 , M_0 at mid span and N_x 0 at support.

- a) Stress Resultants
- i) V_x :— The variation of V_x obtained by classical method is not linear while beam method gives linear distribution. For both the cases magnitude at the bottom of edge beam is greater than that of crown and the ratio of V_x at crown to that at bottom of edge beam for classical method is 0.1765 and for beam method is 0.1222. Again V_x at crown by exact method is 3.287 times greater than that by beam method and at the bottom of edge beam the ratio is 2.805. For the case of classical method magnitude decreases upto the section in between 20°to 10° then bifurcate towards the increasing mode and attains it's maximum value at the bottom of edge beam but for the beam method it follows through out the same path to: attains the maximumat the bottom of edge beam.
- ii) $\mathbb{N}_{\mathbf{x}\emptyset}$:- The mode of the variation of $\mathbb{N}_{\mathbf{x}\emptyset}$ along the depth of the edge beam is approximately similar with a difference that the rate of charge of $\mathbb{N}_{\mathbf{x}\emptyset}$ at the region where it attains it's maximum

of magnitude for the classical theory is smaller than that for the beam theory but the main difference is that the magnitude obtained by the beam method is 50 to 60% less than that obtained by classical method.

- iii) Ng: The mode of the variation of Ngexhibits it's tendency to decrease towards the edge of shell both for classical and beam method. At the crown Ng obtained by the beam theory is 0.1945 times lesser than that obtained by classical theory while at the shell edge it is 2.318 times greater.
- iv) M_{\emptyset} : Sign of M_{\emptyset} , obtained by both the theories is positive. The mode of variation shows the decrease in magnitude as it proceeds from crown and ultimately 0 at the shell edge. For classical theory magnitude decreases upto a section where \emptyset is 30° and then increases upto section $\emptyset = 40^{\circ}$. From this point the magnitude again decreases and comes to 0 at the bottom of the edge beam. Beam theory always gives the lower value of M_{\emptyset} than classical theory.

b) Deformations

The defromations may be calculated by the classical theory from the eqns. of 2.19 by putting V,S,H,M in place of V_2 , S_2 , H_2 and M_2 , where eqn. of 2.26 gives the values of V,S,H and M. The deformations of the example 2 is calculated as follows,

$$\begin{split} \dot{w}_{v} &= \mathbf{c}_{s} & \mathbf{c}_{v1} V + \mathbf{c}_{v2} S + \mathbf{c}_{v3} H \\ &= \frac{1^{4}}{R^{3} t E_{s}} \left[42.38 \times 1744 - 22.99 \times 3568 + 0.9856 (-7477) \right] \\ &= -48 \text{ m.m.} \\ \dot{w}_{h} &= \mathbf{c}_{s} \mathbf{c}_{h1} V + \mathbf{c}_{h2} S + \mathbf{c}_{h3} H \\ &= \frac{1^{4}}{R^{3} t E_{s}} \left[22.99 \times 1744 - 13.59 \times 3658 + 0.2931 (7477) \right] \\ &= -35 \text{ m.m.} \\ \dot{w}_{m} &= \frac{R^{2}}{E_{s} I} \left[0.1245 \times 1744 - 0.07595 \times 3568 + 0.001413 (-7477) \right] \\ &= -0.017 \text{ radians.} \end{split}$$

Deformations by beam theory are calculated as follow:

$$w_v = \frac{4}{\pi} \frac{1^4}{E_s I_s} \left[\frac{1^4}{\pi^4} \left(q_t - \frac{8g P_e}{1^2} \right) - \frac{1^2}{\pi^2} P_e e_e \right]$$
= 3.5 m m

 m_h and m_h can not be determined by beam theory as it neglect those quantities by virtue of first assumptions of article 1.7(a).

Beam theory gives the magnitude of w_v less than that of classical theory, as in this case w_v obtained by classical theory is 13.21 times greater than that determined by beam theory in opposite divection. Beam theory gives w_h and w_m negligible

$$\delta_{\text{vh}} = \frac{(\bar{w}_{h}/L) \text{ obtained by classical theory}}{(\bar{w}_{v}/L) \text{ obtained by classical theory}}$$

$$\delta_{\text{vm}} = \frac{(\bar{w}_{m}/\bar{p}_{k}) \text{ obtained by classical theory}}{(\bar{w}_{v}/L) \text{ obtained by classical theory}}$$

3.4 INFLUENCE OF PRESTRESSING FORCE

of

To investigate the influence presstressing force, shell without prestressing force is considered as reference. The stress distribution for the same shell of example 1. of article 2.6 is shown in fig. 3.1. It is found that the agreement of N_x and N_{x0} obtained by beam and shell theory is close enough but for N_0 and M_0 it differs. The contribution of P_e and external loads along with the dead load of shell towards V_x is shown from fig. 3.2 it is observed that contribution of P_e , calculated by exact theory, is more prominent near the region of shell edge & edge beam and less significant towards crown* while it's

^{*} and it's variation is highly nonlinear along the depth of the transverse section of the shell.

influence follows linear path for beam theory. Both the method give contribution of $P_{\rm e}$ as tensile stress at crown and compressive at bottom. The influence of $P_{\rm e}$, evaluated by beam method, at the crown is 7.89 times greater than that, calculated by exact method while for the bottom of edge beam the ratio is 1.81.

3.5 EFFECT OF BOUNDARY COMDITIONS

The boundary conditions and the deformations of the shell, with and without prestressing force, determined both by beam and classical theory are given in the table 3.1

Table 3.1 DEFORMATIONS OF SHELL EDGE

Boundery	Assumption of beam theory	Deformations				
condition for class-		P _e =400,000 kg		$P_{e} = 0$		
ical theory		Classical	Beam	Classical	Beam	
$w_{vs} = w_{vb}$	wvs wb (assumption 1 of section 1.7(a))	w _v = -48 mm	3.5m.m	21 mm	15 •6	
H = 0 (w _h exists)	w _h =negligible (assumption 2)	w _h =-35 m.m.	negli- gible	- 1 mm	negli- gible	
M = 0 (w _m exists)	w _m = negligible	w _m =-0.077 (rad)	negli- gible		negli - gible	

For $P_e=0$, $\delta_{vn}=0.0476$ and $\delta_{vm}=0$. So for the shell without prestressing force w_h and w_m may be assumed as negligible but with prestressing force they can not be assumed as negligible as it is observed in the previous article that for $P_e=400,000$ $\delta_{vh}=0.728$ and $\delta_{vm}=1.9$. Again for

$$P_e = 0$$
, $\frac{(w_v) \text{ determined by classical theory}}{(w_n) \text{ determined by beam theory}} = 21.0 = 1.345$

while for no prestressing it is 13.21. Therefore the assumptions for boundary conditions in classical and beam theory agree/when $P_{\rm e}=0$ which is not true for $P_{\rm e}=400,000$ kg.

3.6 PARAMITRIC STUDY

The structural behavi-/ of a cylindrical shell is very sensitive to the effect of certain parameters. Therefore it is desirable to study the influence of different parameters such as geometry of the shell, magnitude of prestressing force and eccentricities of the cables on the structural behaviours. In this study the effect of parameters considered are,

- 1. Span to radius ratio, L/R
- 2. Semi central angle, $\mathscr{O}_{\mathbf{k}}$
- 3. Depth of edge beam, d
- 4. Prestressing force, Pe
- 5. Ecentricities of the cable, \mathbf{e}_{o} and \mathbf{e}_{m}

Variation of one parameter is affected at a time while others are kept constants to observe it's influence on the stress resultants and the deformations.

As the arch analysis is dependent on the beam analysis and if beam analysis yields accurate enough longitudinal and shear stress resultants, then My and Ny could be determined approximately correct by proper arch analysis satisfying the boundary conditions. Therefore paramatric study is done for $T_{\rm x}$, $N_{\rm x}$, and defromations. In the figures the variation of $N_{\rm x}$, is shown where longitudinal stress calculated by beam theory is close enough with that determined by classical theory.

The non-dimensionalised parameters which are considered for the study are given in table 3.2

Table 3.2 NON DIMENSIONALISED PARAMETERS

Shell parameters	Non Dimensionalised form	Expression
Semi-central angle	$\mathscr{O}_{f k}$	$ oldsymbol{\emptyset}_{\mathbf{k}} $
Radious, R	Span to radious ratio	I∕ R
Depth of edge beam, d	Depth to span ratio	d∕L
Prestressing force Pe	Coefficient of prestressing force, p	$p = \frac{P_e}{q_{le} I^2}$
Eccentricities	End eccentricity to depth rat	io e _O /d
	Mid eccentricity to depth rat	io e _m /d

$$q_{le} = q_u + q_g \not p_k / \sin \not p_k$$

i) T_x : Longitudinal stress, T_x at the crown decreases as \emptyset_k increases from 30° to 45° in all L/R ratios but increases as L/R increases for all values of \emptyset_k for both p=0.4 and p=0. At the bottom of the edge beam the behaviary is just reverse. Both beam and classical theory yields the above behavi-/ beam theory gives the value of σ_{x} at crown lower than that of classical theory for all values of L/R and \emptyset_k when p=0.4 but for p=0 it always gives higher value than classical theory. At the bottom of the edge beam, beam theory allways underestimates the value of $\tau_{\rm x}$ for p=0.4 but for p=0 it gives higher value than classical theory for L/R=5and $\emptyset_k = 30$, L/R = 10 and $\emptyset_k = 30^\circ$ and 40° . Rest of the cases it gives lower value than the classical theory. The agreement in the variation along the depth of the transverse section does not achieve for p = 0.4. When p = 0 the agreement is there and beam theory gives the varation of $T_{\mathbf{x}}$ accurate enough except the two cases of L/R = 3.33 when \emptyset_k = 40° and 45°. ∇_x attains it's minimum value in between the shell edge and $\emptyset = 10^{\circ}$ for all the cases except for $\phi_k = 30^{\circ}$ and 40° when L/R = 10 and maximum value at the bottom of the edge beam by classical theory when p = 0.4. For the same value of p beam theory gives minimum value at the crown and maximum value at the bottom. T_x is of compressive nature by both the theory for p = 0.4 all along the depth. But for p = 0 maximum compressive stress occurs at the crown and maximum tension occurs at the

bottom of the edge beam by both beam and classical theory.

- ii) $N_{\chi 0}$: For non prestressing stress there is good agreement between beam and classical theory in the distribution of $N_{\chi 0}$. The minimum error is 3.5% while the magnitude of maximum error is 11%. For prestressing shell no comprasion exists in between beam and classical theories. For reinforced shell error increases when smaller values of \emptyset_k has been taken for higher value of I_{χ}/R .
- b) Deformations
- i) Vertical deformations: When p=0.4 the vertical deformation calculated by classical theory, gives the magnitudes much higher than those determined by beam theory. By beam theory difference between the vertical deformations for different values of \emptyset_k is no significant where for classical theory it is opposite. For p=0 mode of the variation is more or less same and the difference of magnitude is not much.
- ii) Horizontal deformations: Both for p=0 and $p=0.4~w_h$ can not be determined by beam theory while classical theory gives significant magnitude for both the cases.
- iii) Rotational deformations: Mode of variation for p=0 is similar with that of $w_{\mathbf{h}}$. Part when p=0.4 it is to some extent different Classical theory gives significant magnitudes for both p=0 and p=0.4 while beam theory neglects it's magnitudes.

The variation of $\mathbf{w_v},~\mathbf{w_h}$ and $\mathbf{w_m}$ with L/R and $\mathbf{\emptyset}_k$ are shown in:fig. 3.10.

3.8 EFFECT OF THE DEPTH OF EDGE BEAM

To study the effect of the edge beam on the structural behaviour the values of d/L ratio consideration are 0.05, 0.04, 0.03, 0.02. Other parameters are kept constant, as given, L/R = 0.03, $\not \!\!\!/_k = 30^\circ$, R/t = 100, b/t = 1.25, p = 0.4 and 0, $e_0/d = e_m/d = 0.4$

a) Stress Resultants

The behavior of the stress resultants are shown in figure 3.5 and 3.6 and the following observations are made:
i) $V_{\rm X}$: When p = 0 the agreement between the beam and the shell theory in the variation is very much dominant for all the values of d/L. When d=0.4 compression occurs both at the top and bottom by classical and beam theory where $V_{\rm X}$ at crown has lower value than at the bottom of the edge beam for a value of d/L 0.5 to 0.4 classical theory gives tension at the edge of the shell at d/L=0.5 and it's magnitude decreases as d/L decreases and the shell is totally in compressive state when d/L=0.3.

Beam theory does not give any where tension. For d/L=0.02 the value of $V_{\rm X}$ at the top/always less than at the bottom and shell is fully in compressive state by both the theories. The agreement in the variation/slightly achieved as d/L decreases.

ii) $N_{\chi 0}$: No agreement between beam and classical theories was found for prestressed shell. For reinforced concrete shell beam theory gives accurate enough stresses which is well comparable with classical theory.

b) Deformations

Variation of deformations with d/L are given in fig.3.10. From the study of this curve following observations are made. i) Vertical deformation (w_y) : For p = 0 at d/L = 0.02 beam theory gives 0.722 time vertical deformation determined by classical theory but at d/L = 0.05 beam theory yields 1.49 times greater defrom tion than the classical theory. When p = 0.4, deformation by beam theory/always less than the that determined by classical theory. As d/L increases defference between the defromation is increased. Classical theory always gives / deformation but for beam theory it is negative for the range d/L = 0.04 to 0.05. W_v decreases for both p = 0 and p = 0.4 and magnitude for p = 0is always greater than the magnitude at p = 0.04 by both the theories. ii) Horizontal deformation (wh): By beam theory horizontal deformation is negligible and it's magnitude can not be determined. determined by classical theory, for p = 0.04/always negative and the magnitude increases as d/L increases. For p = 0 wh is possitive and decreases d/L increases. At d/L=0.0325 it is 0 and for other values it is negative.

iii) Rotational deformation $(w_{\underline{m}})$: The variation of $w_{\underline{m}}$ with d/L is of similar type as that of $w_{\underline{h}}$.

3.9 EFFECT OF ECCENTRICITY:

The following combination of eccentricities are considered to study the effect of eccenticities in structural behavior.

1.
$$e_0/d = 0$$
 2. $e_0/d = 0$ 3. $e_0/d = 0.4$ $e_m/d = 0$ 4. $e_0/d = 0.4$ 5. $e_0/d = 0.2$ 6. $e_0/d = 0.4$ $e_m/d = 0.4$

other parameters are kept constants as follows

L/R = 3.33, d/L = 0.03, R/t = 100, b/t = 1.25, p = 0.4 our From the study of the behavi-/ curve for the stress resultants following observations are made.

i) \mathbb{T}_x : For the combination 3 the agreement between beam and classical method is close. For combination 1 defference between the \mathbb{T}_x determined by beam theory and classical theory is not much. Combination 6 exibits more-or less same manner of variations for combinations 2,4 and 5 no agreement exists. Both the theories keep the shell in the compressive state through out the transverse depth. Beam theory gives lower values of \mathbb{T}_x at the crown and bottom of edge beam than the values determined by classical theory. But \mathbb{T}_x determined in the region of mid depth by classical theory. has lower values than those determined by beam theory. The

influence of eccentricities on \mathbb{T}_{x} at the crown have little while at the bottom of the edge beam they have significant influence.

ii) $N_{\chi 0}$: For all the combinations, except 1, the difference between $N_{\chi 0}$ determined by the classical theory and then beam theory decreases as $e_{\mu \nu}/d$ increases and e_{0}/d decreases, both the theories gives lower values of $N_{\chi 0}$ when $e_{m}/d=0.4$ and for the values of e_{0}/d greater than the e_{m}/d , both the theories gives higher values of $N_{\chi 0}$. Eccentricities have very sensetive influence on the magnitude of $N_{\chi 0}$ for both the theories.

3.10 EFFECT OF PRESTRESSING FORCE:

Values of p considered for this study are 0.8, 0.6 , 0.4, 0.2 and and other parameters are considered as follow: L/R=3.33 $\ell_k=30^{\circ}$, R/t-10.9, b/t=1.25, d/L=0.03.

a) Stress Resultants:

The behaviour curve for stress resultants are shown in Figure 3.8. From the study following observations are made.

i) T_x : For p = 0.8 beam theory give T_x at the crown as tension while for classical theory it is compression for p = 0.2 and 0 both the theories give compression at the top and tension at the bottom. For other three values of p T_x at bottom is compression and the magnitude decreasing with

when p is decreasing. From both the theory it is observed \int_X at crown increases as p decreases. The agreement between beam theory and classical theory is close for p = 0.2 and 0.

ii) N_{xØ};

The differences between the magnitude obtained by beam and classical theory decreases as p decreases and close agreement achieved for p=0.2 and 0.

b) Deformations:

The behaviour curve shown in Figure 3.11.

i) w_v:-

The variation of w_v with p is linear for both the case. Beam theory gives the magnitude always less except when p = 0.2. For p = 0.4 w_v has very small magnitude in posetive direction by beam theory while by classical theory its magnitude is about 8 times greater in opposite direction.

ii) w_h:-

The variation is linear. At p=0 it is magnitude is not 0 but very smalland attains maximum magnitude at p=0.8 by classical theory while beam theory assume it as negligible.

iii) w_m:-

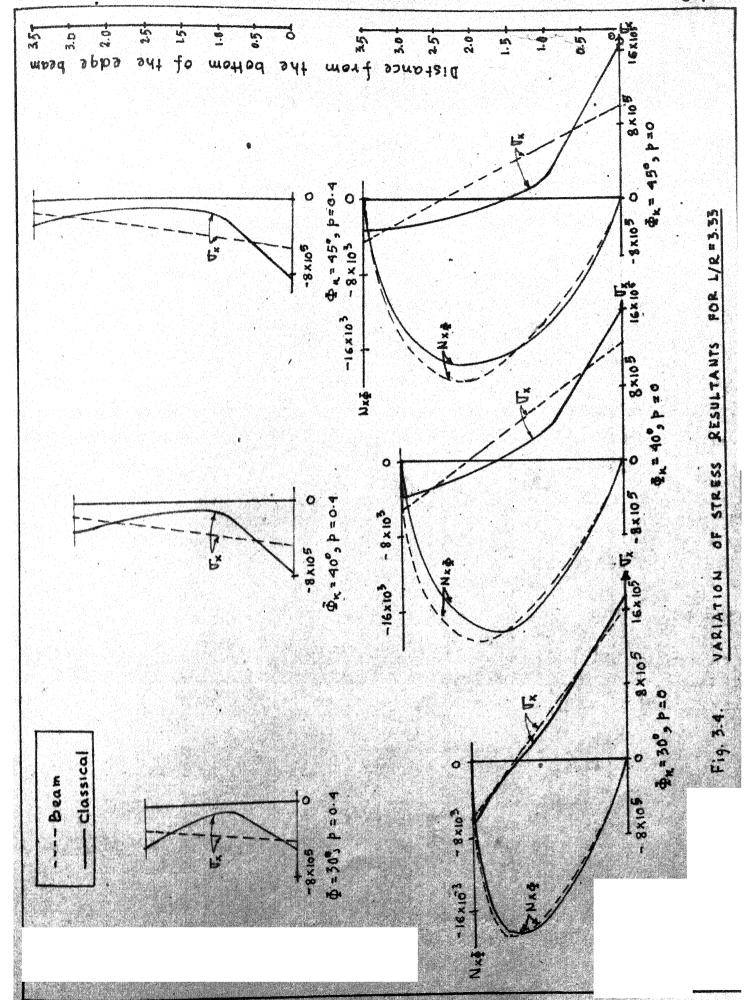
Same type of variation as that of $w_{\mbox{\scriptsize h}}$ only difference is it is magnitude at p=0 is 0.

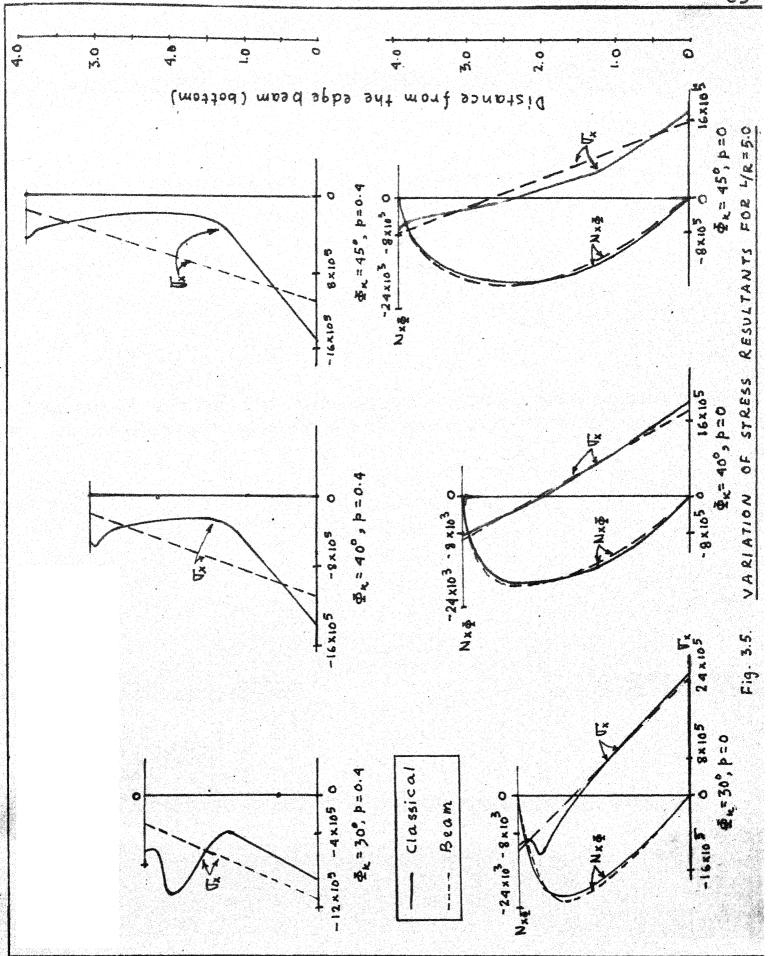
3. DISCUSSION:

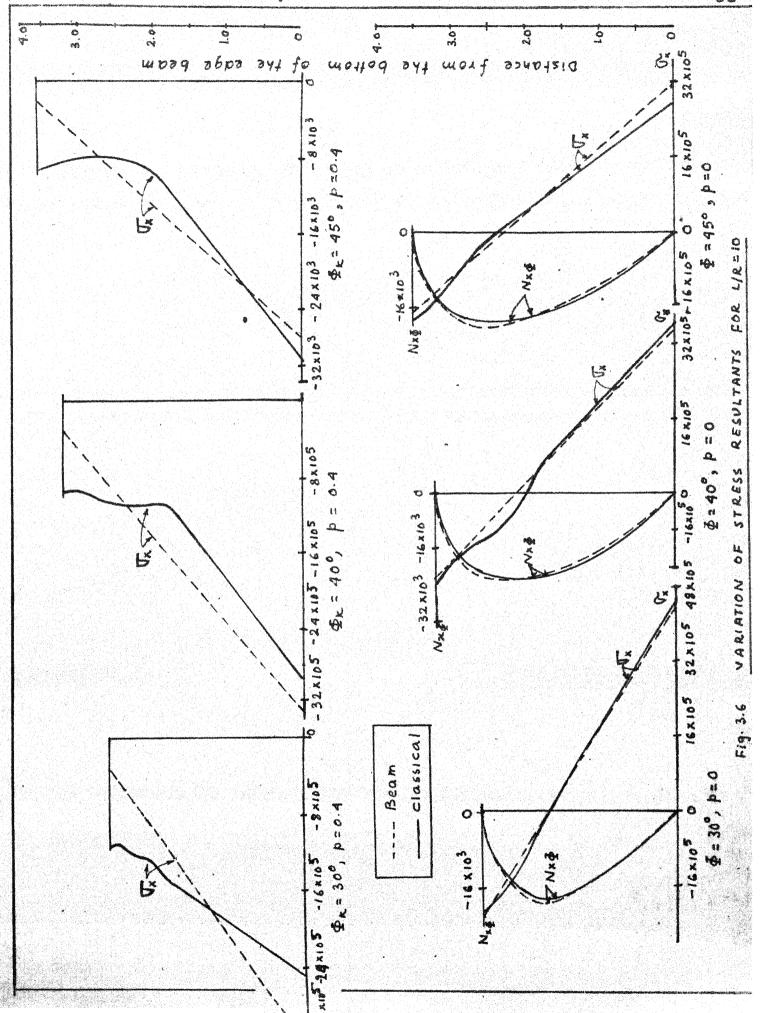
From the above investigations following observations are made,

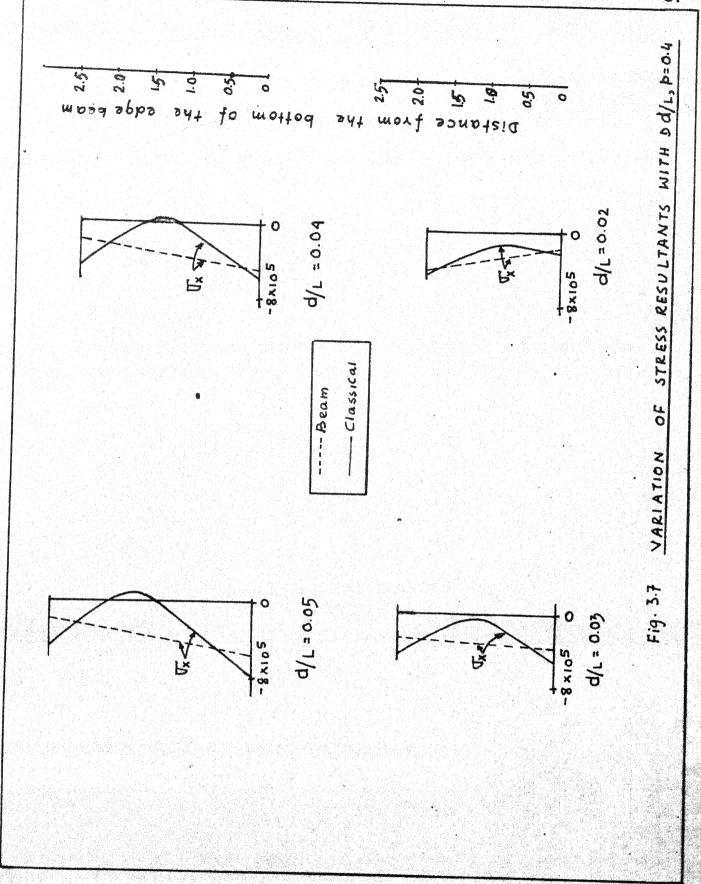
- i) The discrepency between the boundary conditions and the assumptions of beam theory is more significant for prestressed shell.
- ii) Beam theory can be applicable for all values of L/R , d/L and \emptyset_k when p = 0 .
- iii) Beam theory mag be applicable for a limited range of shell parameters for prestressed shell.
- iv) Restraints, fully or partially, yields the stress resultants which compensates each other for $\mathbf{w}_m,\ \mathbf{w}_h$

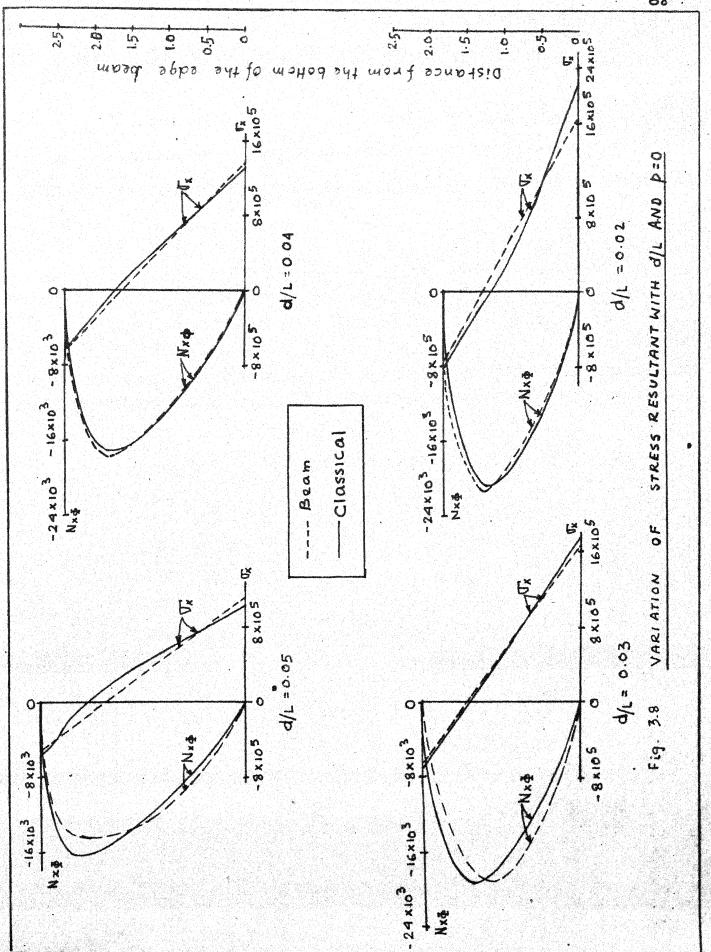
In the next chapter on the basis of the observations, made from the above investigations, theory for arch analysis developed using casting liano's second theorem taking consideration the variation of loads with Sin $\frac{n_\pi x}{1}$.

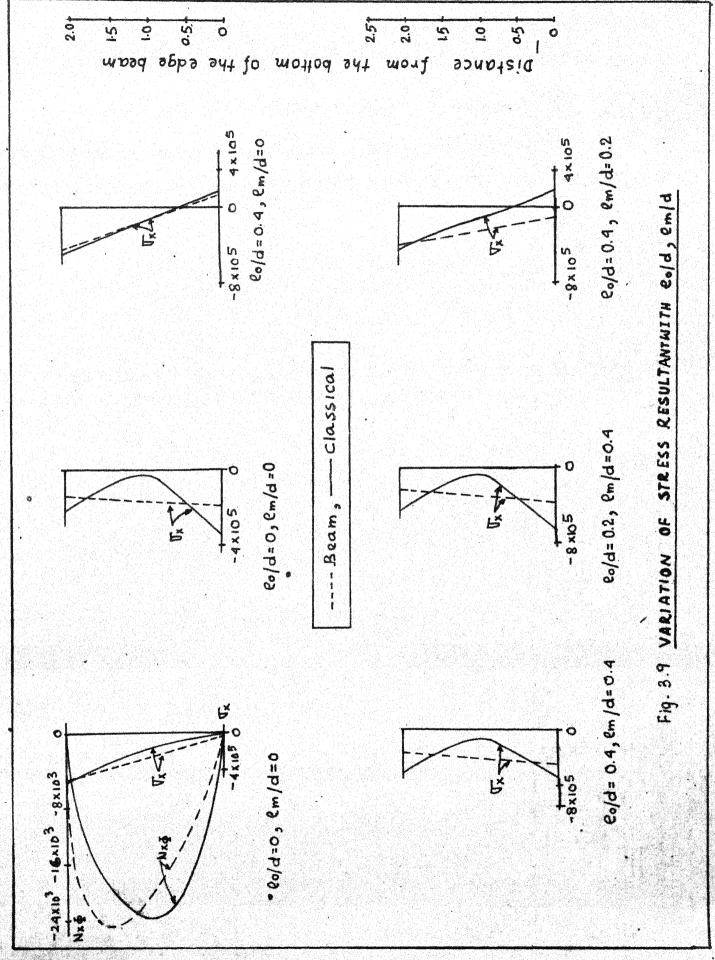


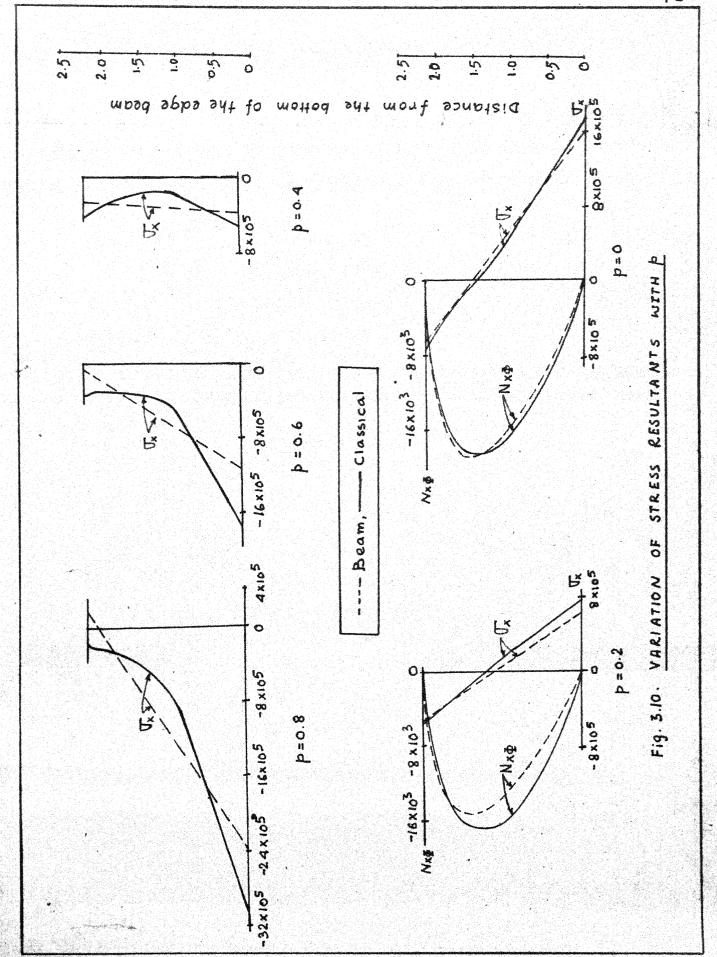


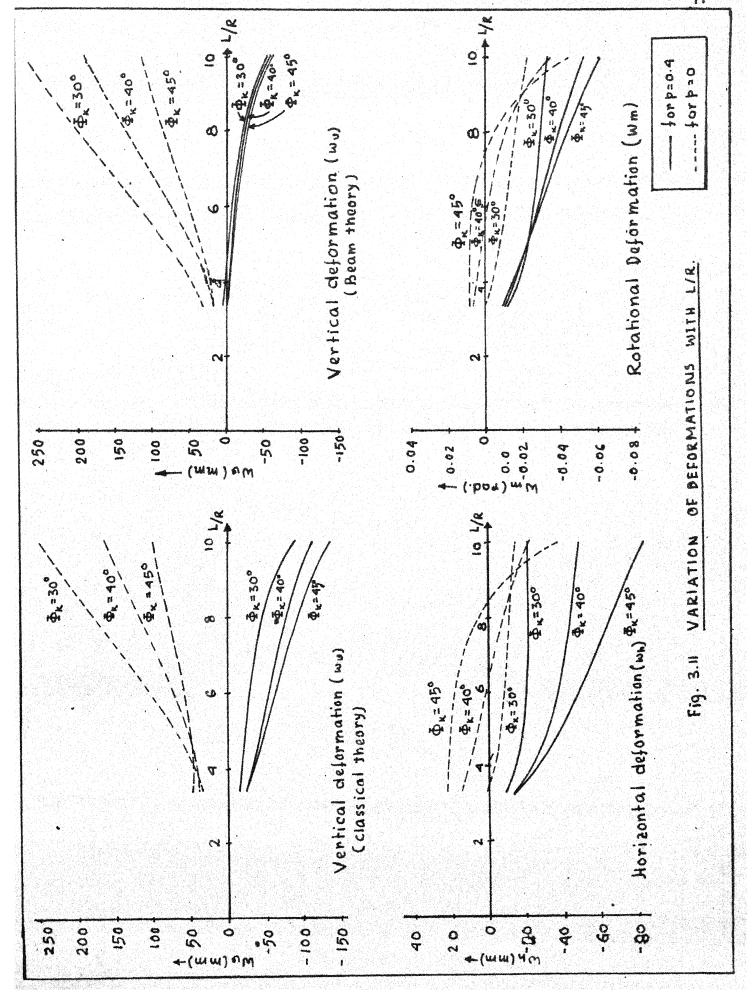




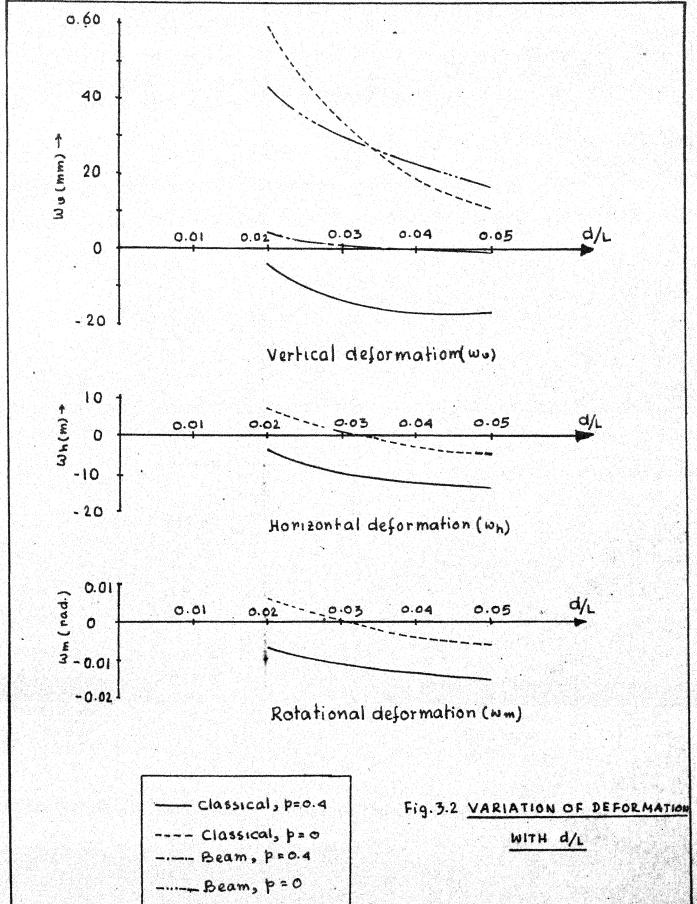


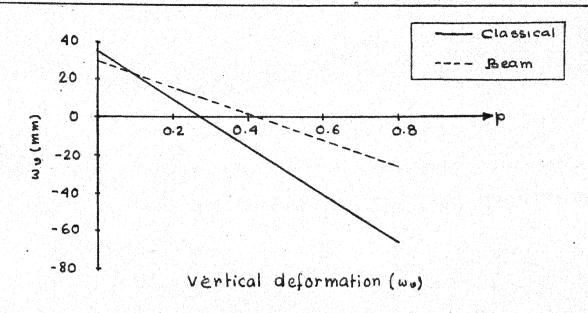


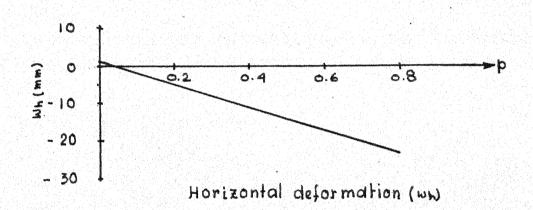












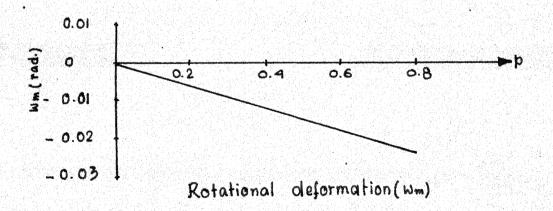


Fig. 3.13 VARIATION OF DEFORMATIONS WITH P

CHAPTER-IV

DESIGN OF CYLINDRICAL SHELL BY SIMPLIFIED BEAM AND ARCH ANALYSIS 4.1 IMPRODUCTION:

It is evident from Figs. 3.3 , 3.4, 3.5 and 3.7 that beam theory yields very good approximate stress distributions for $N_{\rm x}$ and $N_{\rm x}$ for almost all combinations of L/R , $\emptyset_{\rm k}$ and d/L for non prestressed shell. But in article 3.4 it was found that the distributions for $N_{\rm y}$ and $N_{\rm y}$ by beam theory for single barrel shell do not agree with classical theory. Again it is a well established fact that beam theory yields accurate enough stress distributions for $N_{\rm y}$, $N_{\rm y}$ and $N_{\rm y}$ for the interior shell of multibarrel shell group with and without edge beam when p=0. The object of this chapter is to establish such a formulation by which any type of cylindrical shell can be analysed and designed by beam and arch analysis.

Like conventional beam method this method has two distinct parts, beam analysis and arch analysis. The beam analysis does not change but the arch analysis is different. The theory of proposed method is developed assuming the loading on the shell varies with Sin $n \pi x/L$.

In addition to the assumptions made in article 1.7, the method assumes, "The principle of superposition holds, in which case the strain energy is a homogenious function

of second degree, where the strain energy can be expressed as a function of external forces". This assumption is mainly for arch analysis.

4.2 BEAM AMALYSIS

The beam analysis is all most similar to that which was described in article 1.7 (b). The loading is assumed to vary along the longitudinal direction as half sine series.

Loading on shell $q = q_0$ Sin mx

Where
$$q_0 = l_c \left(q_u + \cancel{p}_k \right)_{sin} q_g$$

Moment at any longitudinal section, $M = \frac{q_o}{m^2}$ Sin mx

Shear force at any section, $S_1 = \frac{1}{m} q_0 \cos mx \ c.G.$, area and moment of inertia of, the sections are determined as before, only Q_1 at any node 1 determined as follow,

$$Q_i = 2 R g_i t (Z_s - R k_{zi})$$

At any longitudinal section,
$$T_{x} = \frac{Z_{yt}}{I_{s}} = \frac{1}{m} q_{0} \sin mx$$
At any longitudinal section, $N_{x0} = \frac{Q_{i}}{2} \frac{1}{m} q_{0} \cos mx$

Specific shear, $\frac{\partial N_{x0}}{\partial x} = -\frac{Q_{i}}{2} \frac{1}{L} q_{0} \sin mx$ (4.3)

4.3 ANALYSIS OF CIRCULAR ARCH RESTING ON ELASTIC FOUNDATION

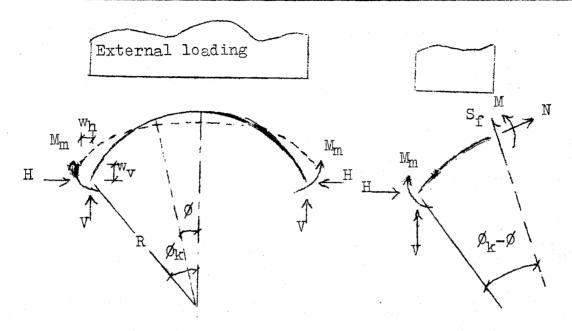


Fig. 4.1 CIRCULAR ARCH ON ELASTIC FOUNDATION

Moment at any section Ø,

$$M = VR(Sin\theta_k - Sin\theta) + HR(Cos\theta_k - Cos\theta) + M_s + M_m \qquad (4.4)$$

$$N = -V \sin \theta - H \cos \theta + N_{S}$$
 (4.5)

$$S_f = -V \cos \emptyset - H \sin \emptyset + S_{fS}$$
 (4.6)

Where, Mg = Moment due to external loading

M = Elastic redundant moment at support

 N_s = Axial force due to external load

Total strain energy of the arch,

$$U = \int_{0}^{S} \frac{1}{2 EI} ds + \int_{0}^{S} \frac{1}{2AE} ds$$

From Castigliano's, second theorem,

Vertical deformation at the support,
$$w_{v1} = \frac{\partial U}{\partial V}$$
 (4.7 a)

Horizontal deformation at the support
$$w_{h1} = \frac{\partial U}{\partial H}$$
 (4.7 b)

Rotational deformation at the support,
$$w_{m1} = \frac{\partial U}{\partial M_m}$$
 (4.7 c)

Futting eqns. 4.4 and 4.5 on eqs. of 4.6, and equating w_{v1} , w_{h1} , w_{m1} to w_{v} , w_{h} and w_{m} ,

$$(T^3k_{mv} + Tk_{mv}) V + (T^3k_{mvh}) H + (T^3k_{mvm}) M' = Ew_v - T^3k_{11}$$
 (4.8)

$$(T^3k_{mvh}) V + (T^3k_{mh} + Tk_{nh}) H + (T^3k_{mhm}) M' = Ew_h - T^3K_{21} + TK_{22}(4.9)$$

$$(T^3k_{mvm}) V + (T^3k_{minm}) H + (T^3k_{mm}) M' = EV_m - T^3K_{31}$$
 (4.10)

Where,

$$k_{mv} = 12 p_k (2 \sin^2 p_k + 1) - 6 \sin 2 p_k$$
 (4.11 a)

$$k_{mh} = 12 p_{k} (2 \cos^{2} p_{k} + 1) - 18 \sin 2 p_{k}$$
 (4.11 b)

$$k_{mm} = 24 p/k$$
 (4.11 c)

$$k_{mvh} = 12 \beta_k \sin 2 \beta_k - 24 \sin^2 \beta_k \qquad (4.11 d)$$

$$k_{mvm} = 24 p_k \sin p_k \qquad (4.11 e)$$

$$k_{\text{mhm}} = 24 \%_{k} \cos \%_{k} - 24 \sin \%_{k}$$
 (4.11 f)

$$k_{mv} = \phi_k -0.5 \sin 2\phi_k \qquad (4.11 g)$$

 $M \cdot = M_m/R$

$$k_{\rm nh} = p_{\rm k} + 0.5 \sin 2p_{\rm k}$$
 (4.11 h)

$$K_{11} = \int_{-\phi_{k}}^{\phi_{k}} \frac{M_{s}}{R}$$
 Sinf. as (4.12 a)

$$K_{21} = \int_{-\emptyset_{\mathbf{k}}} \frac{\mathbb{M}_{\mathbf{s}} \left(\cos \emptyset_{\mathbf{k}} - \cos \emptyset \right) \, \mathrm{d}\emptyset}{\mathbb{R}}$$
 (4.12 b)

$$K_{22} = \int_{\Gamma} \int_{\mathbf{k}} \cos \theta \, d\theta \qquad (4.12 c)$$

$$K_{31} = \int_{-\phi_{\mathbf{k}}}^{\phi_{\mathbf{k}}} \frac{M_{\mathbf{s}}}{R} d\phi \qquad (4.12 d)$$

$$T = \frac{1}{a} \frac{R}{t} \tag{4.13 a}$$

$$T^3 = \frac{1}{a} \left(\frac{R}{t}\right)^3$$
 (4.13 b)

a = width of the radial section of the arch

t = depth of the radial section of the arch

Eqns. 4.7, 4.8, and 4.9 have 3 un knowns, V, H and M and hence will give unique solution to determine the values of V,H,M, and putting those values on eqns. 4.4, 4.5 and 4.6 M, N and S_f can be determined.

To derive the eqns. 4.8, 4.9, 4.10 the contribution of $S_{\hat{\Gamma}}$ to strain energy is neglected and the radial sectional properties of the arch are assumed to be constant with the variation of δ . The external loading is also taken symmetric about Z-axis.

4.4 ARCH A MAIYS IS

For the arch analysis for the proposed method the elementary strip of the circular arch is considered and this elementary arch strip is considered to be supported on the edge beam. This support is assumed to be elastic support which can deforme in three directions, i.e., vertical, horizontal and rotational directions. In general each support point has three degrees of freedome while actual conditions are determined from the boundary conditions of the shell. The support reactions are determined by satisfying the compatibility of the deformations of shell edge and edge beam at the interface in mid span.

The elementary arch strip is analysed by the method described in article 4.3. The edge beam is assumed to offer restraint against vertical deformation only. Again considering the statical equilibrium the arch is in equilibrium under the action of external loading and specific shear.

As the arch strip in separated, for arch analysis, from edge beam then it would be in equilibrium under the loading of $F_{\mathbf{v}}$ and V

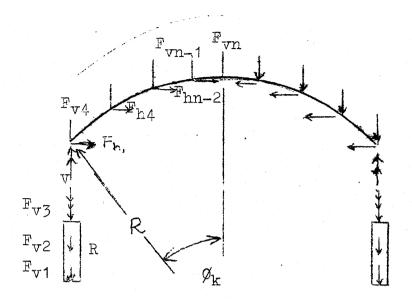


Fig. 4.2 STATICAL EQUILIBRIUM OF ELEMENTARY ARCH STRIP

Therefore, the arch analysis essentially stands as the solution of two simultaneous eqns. 4.9 and 4.10 satisfying the conditions given by eqn. 4.14. Hence solution may be achieved solving two simultaneous Equations. Putting the the value of V from eqns. 4.14 to 4.9 and 4.10

$$(\mathbf{T}^{3}\mathbf{k}_{\mathrm{mh}} + \mathbf{T}\mathbf{k}_{\mathrm{nh}}) + (\mathbf{T}^{3}\mathbf{k}_{\mathrm{mhm}}) = \mathbf{E}_{\mathbf{S}^{W}\mathbf{v}} - \mathbf{T}^{3}\mathbf{k}_{21} + \mathbf{T}\mathbf{k}_{22} - (\mathbf{T}^{3}\mathbf{k}_{\mathrm{mvh}}) \mathbf{v}$$

$$(4.15)$$

$$(T^3 k_{mhm}) H + (T^3 k_{mm}) M' = E_s W_m - T^3 K_{31} - (T^3 k_{mvm}) V$$
 (4.16)

Where, T = R/t

The unique solution of eqns. 4.15 and 4.16 will determined the values of H and M. As soon as V,H,M are known Mg and Ng at any longitudinal section, are given by:

$$\mathbb{N}_{S} = \mathbb{N}_{S} + \mathbb{R} \mathbf{c}_{\mathbf{n} \mathbf{V}} + \mathbb{R} \mathbf{c}_{\mathbf{n} \mathbf{H}} + \mathbb{R} \mathbb{N}^{*}, \qquad (4.17)$$

$$\mathbb{N}_{g} = \mathbb{N}_{s} + c_{nv} \mathbf{V} + c_{nh} \mathbf{H}$$
 (4.18)

where
$$c_{mv} = \sin \phi_k - \sin \phi$$
 (4.19 a)

$$\mathbf{c}_{mh} = \cos \phi_{k} - \cos \phi \tag{4.19 b}$$

$$c_{nv} = -\sin \theta \tag{4.19 c}$$

$$e_{nh} = -\cos \emptyset \tag{4.19 d}$$

4.5 BOUNDARY CONDITIONS

For single barrel shell H=0 and M'=0 as it is assumed that edge beam does not offer any restraints towards horizontal and rotational deformations.

For the interior shell of multibarrel group $\mathbf{w}_h = \mathbf{0}$ and $\mathbf{w}_m = \mathbf{0}$.

4.6 SUMMARY OF COMPUTATIONS

For the systematic and easy computional work of cylindrical shell analysis, the proposed method is detailed below in 6 steps.

- 1. Sectional Properties (Ref.Fig.2.3)
- a) Assume geometrical properties of the transverse section.

b)
$$h = R (1 - \cos \theta_k) + d$$

c)
$$l_c = 2R \sin \theta_k$$

d) Area of cross section,
$$A_c = 2 \text{ Rt} \emptyset_k$$

$$A_b = 2 bd$$

$$A_s = A_c + A_b$$

$$z_c = R k_z$$

$$Z_b = h - d/2$$

$$Z_{s} = (A_{c}Z_{c} + A_{b}Z_{b}) / A_{s}$$

f) Moment of inertia,

$$I_e = R^3 t k_I$$

$$I_h = bd^3/6$$

$$Z_{rc} = Z_{s} - Z_{c}$$

$$Z_{rb} = Z_{b} - Z_{s}$$

$$I_s = I_b + I_c + A_b (Z_{rb})^2 + A_c (Z_{rc})^2$$

g) Co-ordinates of nodal points (Fef. Fig.2.4),

$$d\emptyset = \emptyset_k/k$$

k is always even number such that $d\emptyset > 5^{\circ}$

Number of nodal points n = m+3, m=k+1

$$Z_i = R(1-\cos\phi_i)-Z_s$$

$$Y_i = R \sin \theta_i$$

- h) Distance of top fiber from C.G., $Z_{yt} = Z_{s}$ Distance of bottom fiber from C.G., $Z_{yb} = h-Z_{s}$
- i) First moment of area, i) For shell portion

$$Q_i = 2 Rt (Z_s - R k_{Zi})$$

ii) For beam portion

$$Q_{i} = b (Z_{vb}^{2} - Z_{i}^{2})$$

2. Loading (mef. Fig. 1.1c)

Determine qu and qg

$$q_o = 4 \left[\left(\frac{p_k}{\sin p_k} (q_g + t) + q_u \right) + 1_c + 2bd \right]$$

- 3. Beam Analysis
- a) $T_{xt} = -q_0 (1/x)^2 (Z_{yt}/I_s)$ Sinmx

$$\sqrt{xb} = q_0 (1/\pi)^2 (Z_{yb}/I_b) \sin mx$$

- b) $N_{x/i} = q_0 (1/\pi) (1/2 I_s) Q_i Cosmx$
- c) Specific shear = $(q_0/2 I_s) \not p_i$ Sin mx
- 4. The above computation canbe done systematically by the following tabular form.

	Ø,	Co-ordinate	SS			
Nodal in point degr	ree rad.	z _i Y _i		Q _i 1	xø T	Specific shear

5. Arch Analysis

a) Shell Beam
$$ds_4 = ds_m = Rd\theta/2$$
, $ds_1 = Rd\theta$ $ds_1 = ds_3 = d/4$, $ds_2 = d/2$ $dy_1 = R (sin\theta_{1-1} - sin\theta_1)$ $dy_1 = 0$ $dz_1 = R (cos\theta_1 - cos\theta_{1-1})$ $dz_1 = ds_1$ where, $\theta_4 = \theta_k$, $\theta_5 = \theta_k - d\theta/2$ $\theta_1 = \theta_{1-1} - d\theta$, $i = 6, 7 - - - n - 2$ $\theta_{n-1} = d\theta/2$, $\theta_n = 0$ b) $F_{vi} = q_g dy_1 + \frac{\partial F_{vi}}{\partial x} dz_1$ $F_{hi} = \frac{\partial F_{vj}(Y_j - Y_i)}{\partial x} + F_{hj} (Z_j - Z_j)$ $F_{vj}(Y_j - Y_i) + F_{hj} (Z_j - Z_j)$ $f_{vj}(Y_j - Y_i) + f_{hj} (Z_j - Z_j)$ $f_{vj}(Y_j - Y_i) + f_{hj} (Z_j - Z_j)$

c) The arch analysis should be carried on with the help of above expressioniin the following tabular form.

		U IV				
Nodal ds points	dy dz	$\frac{\mathbf{x}\mathbf{x}}{\mathbf{x}}$ d	z qody	$\mathbf{F}_{\mathbf{v}}$	F _h M	e N

d)
$$V = \Sigma F_{vi}$$

$$T = R/t$$

$$A_{11} = T^{3} k_{mh} + T k_{nh}$$

$$A_{12} = T^{3} k_{nhm}$$

$$A_{21} = A_{12}$$

$$A_{22} = T^{3} k_{mm}$$

$$K_{21} = 12 \int_{-\emptyset_{K}} \frac{M_{S}}{R} (\cos \emptyset_{K} - \cos \emptyset) d\emptyset$$

$$K_{22} = \int_{-\emptyset_{K}} K_{S} \cos \emptyset d\emptyset$$

$$K_{31} = 12 \int_{-\emptyset_{K}} \frac{M_{S}}{R} d\emptyset$$

The above integral may be evaluated by Simpson's rule.

$$C_{1} = E_{s} W_{h} - T^{3} K_{21} + TK_{22} - T^{3} k_{mvh} V$$

$$C_{2} = E_{s} W_{m} - T^{3} K_{31} - T^{3} k_{mvm} V$$

$$A = \begin{bmatrix} H \\ C_{2} \end{bmatrix}$$

$$C_{1} C_{2}$$

$$(4.20)$$

H and M^{\prime} is solved putting boundary conditions on 4.20.

6. Final values of
$$N_{\text{p}}$$
, N_{p}

$$M_{\text{p}} = M_{\text{s}} + Rc_{\text{LW}}V + Rc_{\text{rh}}H + RM'$$

$$N_{\text{p}} = N_{\text{s}} + c_{\text{rw}}V + c_{\text{rh}}H$$

The values of k_{I} , k_{X} , k_{mv} , k_{mh} , k_{mm} , k_{mvh} , k_{mvm} , k_{mhm} , k_{mh} , k_{nh} are given in appendix for different values of \emptyset and \emptyset_k .

4.7 ILLUSTRATED EXAMPLE

Example - 2

Span = 30m,
$$\emptyset_k$$
 = 30°, R = 9m, d=90cm, b =11.25 cm, t=9cm.
Loading: Lime terracine = 45 kg/m

Live load = 75 kg/m²

Total = 120 kg/m²

Step-1

Sectional properties:

$$h = 2.106 \text{ m}, l_c = 9.0 \text{ m}, A_s = 1.0507 \text{ m}^2, Z_s = 0.6466 \text{ m}$$

 $I_s = 0.3795 \text{ m}^4$

 $d p = 5^{\circ}$ where k = 6, n = 7, n = 10

 $Z_{yt} = 0.6466 \text{ m}, Z_{yb} = 1.4594 \text{ m}$

In shell portion $Q_i = 1.62 (0.6466 - 9 k_z)$

In beam portion $Q_i = 0.1125 (2.12 - Z_i^2)$

Step-2

$$q_0 = \frac{4}{\pi} \left(\frac{0.5236}{0.5} \text{ x} 3652.656 + 2 \text{ x} 0.9 \text{ x} 0.1125 \text{x} 2400 \right) = 4650.7061 \text{ kg/m}$$

Step-3

Longitudinal stress at top fiber, $T_{xt} = -717000$ kg/m

Longitudinal stress at bottom fiber, Txb= 1620000 kg/m

 $N_{xp_i} = 58475.1435 Q_i$

Specific shear = $6127.4165 Q_{i}$

Step-4

Table 4.1 STRESS RESULTANTS (x, Nxø

Noada point		Ø In		rdinates Y.	Q (m ³)	(kg/m²)	N xø (kg/m)	Specific shear
	degr	ee radian	Z _i (m)	Y (m)	(111)		(rg/m)	(kg/m²)
1		. 1	1 45 94	4 . 5	0	1620000	0	C
2			1.0094	4.5	0.1250		7310.9125	765.12
3			0.5594	4.5	0.2044		11950.8231	1251 •79
4	30	0.5236	0.5594	4.5	0.2044		11950.1215	1251.79
5	25	0.4363	0.1966	3.8036	0.2571		15020-3541	1574 •53
6	20	0.3490	-0.1018	3.0782	0.2629		15380.3618	1610.05
7	15	0.2618	-0.3399	2.3294	0.2308		13500.7451	1413 -46
8	10	0.1745	-0.5099	1.5629	0.1699		9720-4175	1040.50
9	5 .	0.0872	-0.6123	0.7844	0.0897		5250.3950	549.95
10	0	0	-0.6466	0	0	-717000	0	0

tep - 5

Table 4.2 CALCULATION OF N AND N

M S S	0 0 0	0 505:8734 180:5531 -1298.3477 389.8418 -2453.9975	- 742 • 7895 - 3681 • 4553 - 1152 • 2685 - 4503 • 4460 - 1598 • 1713 - 4914 • 8235	-2019.2314 -4933.6206
d d	000	68.4679 430.9913 87.4407 1120.44 96.8099 1187.8999 –	1071 • 9752 - 761 • 6494 - 430 • 1746 -	î o
H.	77.235 -189.7621 -206.9603	- 68.4679 430.99 -187.4407 1120.44 - 96.8099 1187.89	48.2734 194.1696 299.3735	167.9952
qudy	000	000	000	0
gggs	77.235 155.45 77.235	167.9952 335.9904 335.9904	335 <u>.</u> 9904 335 <u>.</u> 9904 335 <u>.</u> 9904	167.9952
$\frac{\partial N_{\mathbf{x}} \beta}{\partial \mathbf{x}} d\mathbf{z}$	0 -345,2121 -284,1953	-236.4631 -523.4311 -432.8003	0.2032 <u>-</u> 286.7170 0.1363 -141.8208 0.0684 - 36.6169	0
ර වැ.	0.425 0.45 0.225	0.1839 0.3318 0.2685	0.2032 0.1363 0.0684	0,0
dy.	000	0.3927 0.3443 0.7854 0.7116 0.7854 0.7378	0.7854 0.7584 0.7854 0.7732 0.7854 0.7822	0.3915
a d	0.225 0.45 0.225	0.3927 0.3443 0.7854 0.7116 0.7854 0.7378	0.7854 0.7584 0.7854 0.7732 0.7854 0.7822	0.3927 0.3915
Nodal	- a n	4 ΓΛ Φ	► ® O	0

$$V = \sum_{j=4}^{10} F_{vj} = 356.0932$$

The shell is single barrel, so, H = 0 and $M^{i} = 0$ Step - 6

Final stress resultant is calculated and given in table 4.3 and companed with the result calculated by classical theory.

Table 4.3 Mg AND Ng

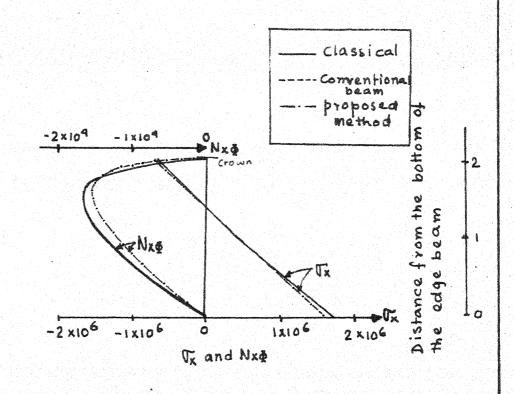
ø c _{mv} _		$^{ m M}\!$		c	$^{ m N}\!$		
	шV	1	2	c m	1	2	
30	0	Ø	0	- 0.5	- 484	- 180	
20	0.15798	121	137	-0.34202	- 2583	- 22 7 0	
10	0.32635	-110	- 96	-0.17365	- 4565	- 425 7	
0	0,5	- 424	- 460	0	- 4934	- 50 1 0	

Column 1 and 2 of M_{\odot} and N_{\odot} stands for the values calculated by proposed method and classical theory respectivty. The variation of stress resultants are shown in figure 4.3.

4.8 DESIGNOOF REINFORGEMENTS

a) Theoretical Aspects

The 4 main stress resultants, for the purpose of design



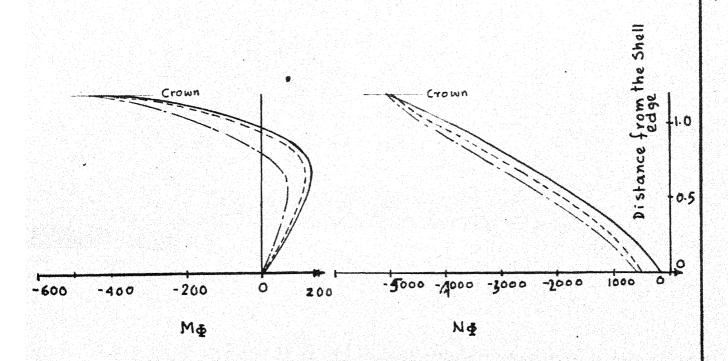


Fig. 43. VARIATION OF STRESS RESULTANTS

of a cylindrical shell roof, are \mathbb{N}_{x} at x=0, \mathbb{N}_{x} , \mathbb{N}_{0} and \mathbb{N}_{0} at x=L/2. On the basis of these stress resultants the reinforcements are devided into two groups, 1) Force reinforcements, 2) Moment reinforcements. These are explained below.

1) Forces Reinforcements

- i) Longitudinal reinforcements, this is being to overcome $\mathbb{N}_{\mathbf{x}}$.
- ii) This is being due to $N_{oldsymbol{\emptyset}}$
- iii) Diagonal tension created by the combination of Ng, Nx and $^{N}\!\mathbf{x}\mathbf{\emptyset}$

 $N_{\rm x}$, $N_{\rm p}$ and $N_{\rm x}$, action on a small element of the shell, may be considered as orthogonal stress in a plane stress system.

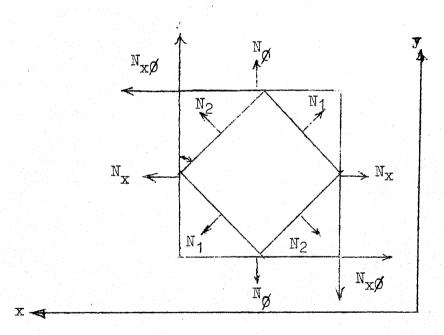


Fig. 4.3 PRINCIPAL STRESSES IN SHELL ELEMENTS

In true theoretical sense the combinationshould be of N_x , N_{g} , N_{xg} , where $N_{xg}' = N_{xg} - M_{xg}/R$. Normally M_{xg} is negligible

Principal stresses,
$$N_{1,2} = \frac{N_x + N_y}{2} \pm \frac{1}{2} \left\{ (N_x - N_y)^2 + 4 N_{xy}^2 \right\}^{\frac{1}{2}}$$
 and the direction is given by $\tan 2\theta = \frac{2 N_{xy}}{N_x - N_y}$ (4.21 a)

Thus eqs. of 4.21 provide all the necessary information, to design the reinforcement, these should be examined at x = 0, L/4 and L/2.

At $x = \mathbf{N}_2$, $\mathbf{N}_{\mathbf{X}} = 0$, therefore $\mathbf{N}_1 = \mathbf{N}_{\mathbf{X}}$ and $\mathbf{N}_2 = \mathbf{N}_{\mathbf{X}}$ where $\theta = 0$ ° or 90°

At
$$x = L/4$$
, $N_{1,2} = \frac{1}{V_2} \left[\frac{N_x = N_0}{2} \pm \left\{ (N_x - N_0)^2 + 4 N_x 0 \right\}^{\frac{1}{2}} \right]$ and $\tan 2\theta = 2 N_x 0 / (N_x - N_0)$.

At $x = \mathbf{D}$, \mathbb{N}_x and \mathbb{N}_y are zero. Thus $\mathbf{N}_{1,2} = \pm \mathbb{N}_x \emptyset$ and $\tan 2\theta = 2 \mathbb{N}_x \emptyset 0 \quad \infty \quad \theta = 45^\circ$.

2) Moment Reinforcements. This is transverse reinforcements to overcome the bending moment M_{\emptyset} .

The analysis of shell has been based on the assumption of homogeneity of the material. In other words, the elastic

20 m.m. \emptyset is provided from bottom to shell upto $\emptyset = 20^{\circ}$.

ii) For Diagonal tension

Maximum $N_{x0} = 15380$ kg/m. Area of steel required is $10.9 \text{ cm}^2/\text{m}$. Provide 12 0 0 10 cm c/c.

iii) For Transverse Force, Nø

 $N_{\rm p}$ is always compressive and maximum value is 5010 kg/m. So no reinforcement is required.

2) Moment Reinforcements

Moment reinforcement is required only for $^{\text{IM}}$ \emptyset and it charges it's Sign at $\emptyset = 14^{\circ}$.

Maximum $M_{\odot} = -424$ Kgm/m. To resists this moment by a singly reinforced concrete section the eff. depth required/5.49 cm If the eff. cover is kept 2 cm then eff. depth available is 7 cm

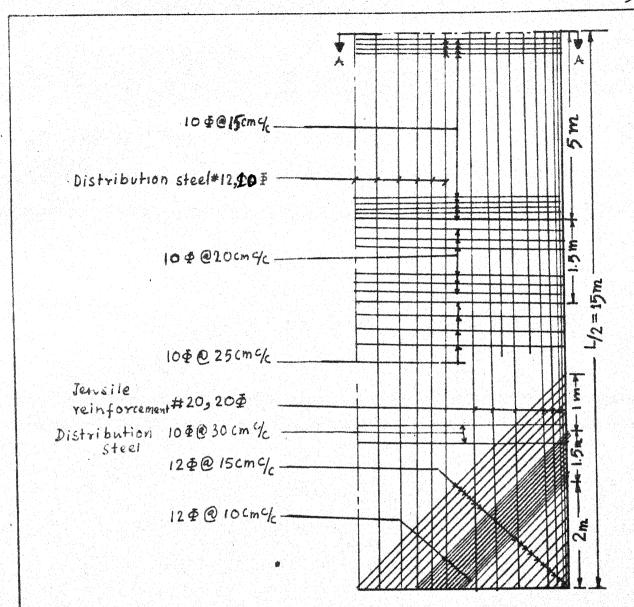
· 0.K.

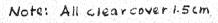
Steel required to resists this mount = $4.97 \text{ cm}^2/\text{m}$, provide 10 \emptyset @ 15 cm c/c.

Distribution steel: 10 Ø . @ 30 cm c/c.

4.9 DISCUSSION

a) The formulations, which have been derived in article 4.4 and 4.5 and presented in 6 steps of calculations are fairly simple enough and can be concieved without the knowledge of shell theory, with the help of the tables given, the computational work is simple and





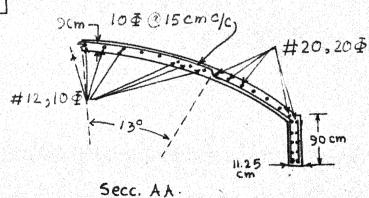


Fig. 4.4 REINFORCEMENT DETAILS (Ex.2)

less tedious. The tables which are given can be used for any values of R/t, L/R and $\theta_{\rm k}$. The single barel, interior and eaterion shell of a multibarrel group can be analysed uniquely Stress resultant can be obtained at any longitudinal section can with $\sin^{n\pi x}$ by these formulations. The error involvement with the calculations of arch analysis is less than that which are done by column analogy method or elastic centre method. The error involvement, with the problem of example -1 of article 4.7, for $T_{\rm k}$ at top and bottom are 5.5% (more) and 6.25% (less), for $T_{\rm k}$ Mg and $T_{\rm k}$ errors are 15% (less), 8.5% (more) and 1.5% (less) respectivly for maximum values

b) The expression, of step 5(d) of the article 4.6 has been derived satisfying the static equilibrium. This could be done by satisfying the vertical deformation compatibility. The vertical deformation of the edge beam at any longitudinal section is given by,

$$w_{vb} = \frac{1}{E_b} \frac{1^4}{ba.3} (0.12319 \text{ v} -0.19351 \frac{d}{L} \text{ N}_{x} \text{w}_4^+ \frac{4}{\pi} 0.12319 \text{ bd } \text{V})$$

$$\sin \frac{n\pi x}{l} (4.21)$$

From the eqn. 4.8 the vertical deformation of the shell is given as (assuming external load varies with Sin mx)

$$w_{vs} = \frac{1}{E_s} \left[(T^3 k_{mv} + Tk_{nv})^2 W_+ (T^3 k_{mvh}) H + (T^3 k_{mvm}) M' + T^3 K_{11} \right] \sin Mx$$
(4.22)

To satisfy the vertical deformation compatibility at the shell beam interface $\mathbf{w_{vs}}$ should be equal to $\mathbf{w_{vb}}$. Therefore vertical deformation compatibility yields,

$$(T^{3}k_{mv}+Tk_{nv}-0.12319 c_{b}) V+(T^{3}k_{mvh})H+(T^{3}k_{mvm})M'=1-T^{3}K_{11}+c_{b}(w_{vbg}-w_{vbs})$$
(4.23)

Where,
$$c_b = \frac{E_b}{E_s} = \frac{1^4}{bd^3}$$
, $w_{vbg} = \frac{4}{\pi}$ 0.12319 $bd \checkmark$, $w_{vbs} = 0.19351 N_x / 4 \frac{d}{1}$

The steps 5(d) of article 4.6 will be changed as follow:

Step 5(d) of article 4.6:

$$T = R/t$$
, $c_b = \frac{E_b}{E_s} \frac{1^4}{bd^3}$, $w_{vbg} = \frac{4}{\pi}$ 0.12319 bd , $w_{vs} = 0.19351 \frac{d}{L} N_{xp}$
 $A_{11} = T^3 k_{mv} + Tk_{mv} = 0.12319 c_b$
 $A_{12} = A_{21} = T^3 k_{mv}$

$$A_{13} = A_{31} = T^3 k_{mvm}$$

$$\Delta_{22} = T^3 k_{mh} + Tk_{nh}$$

$$A_{23} = A_{32} = T^3 k_{mhm}$$

$$A_{33} = T^3 k_{mm}$$
 $K_{11} = 12 \int_{K_{k}}^{K_{k}} \frac{M_{s}}{R} \sin \phi_{k} d\phi$ and K_{21} , K_{22} and K_{31} are given

in step 5(d). of article 4.6

Above integral may be evaluated by Simpson's rule.

$$c_1 = T^3 K_{11} + c_b (w_{vbg} - w_{vbs})$$
 $c_2 = E_s w_h - T^3 K_{21} + K_{22}$
 $c_3 = E_s w_m - T^3 K_{31}$

Therefore vertical, horizontal and rotational deformation compatibility at the interface yields,

$$\begin{bmatrix} \mathbf{A} & \begin{bmatrix} \mathbf{V} & \mathbf{C}_1 \\ \mathbf{H} & \mathbf{C}_2 \\ \mathbf{C}_3 \end{bmatrix} \tag{4.24}$$

The boundary conditions are:

- i) For single barrel shell H=0 and M = 0
- ii) For interior shell of multibarrel group, $w_h=0$, $w_m=0$ substituting these boundary conditions m eq. 4.2 and solving them V, H and M' could be defined uniquly. The final stress resultants will be obtained from step 6 of article 4.5.

CHAPTER-V

CONCLUSION AND RECOMMENDATION FOR FURTHER WORK

The following conclusions are drawn based on the present investigation.

5.1 EFFECT OF SHELL DIMENSIONS:

1. Reinforced concrete shells: Magnitudes of stress resultants decrease with increase of semicentral angle while increase with span to radius and depth to span ratios by both beam and classical theories. The influence of membrane solution on the total solution increases with semicentral angle. Beam theory agrees closely with classical theory for any values of radius to thickness and depth of edge beam to span ratios when span to radius ratio is greater than 3.

Accuracy of beam theory depends on semicentral angle to a great extent. Higher value of semicentral angle for higher value of L/R ratio, yields the stress resultants with 4 to 10 % error. The preferable range of depth of edge beam to span ratio is 0.03 to 0.05.

Vertical, horizontal and rotational deformations increase with span to radius and depth of edge beam to span ratios while decrease with the increase of semicentral angle for classical theory. Beam theory also exhibits the same feature of vertical deformations. Restraints of horizontal

and rotational deformations produce stress resultants which are in a nature to compensate each other. The ratio of the magnitude of vertical deformation, determined by classical theory to that determined by beam theory is in the range of 0.262 to 1.43. Depth of edge beam to span does not have much influence as span to radius, ratio and semicentral angle have on this range. The range could be reduced as 0.8 to 1.43. With the proper selection of \emptyset_R and L/R ratio.

by prestressing force for suitable depth of edge beam for all values of span to radius, ratio and semicentral angle by classical theory. Tension occurs at shell edge for depth to span ratio 0.05 and 0.04. Beam theory does not give any tension for higher values of depth to span ratio. No exact conclusion could be drawn about the influence of semicentral angle and span to radius ratio on the agreement between classical and beam theories in the variation of stress resultants. No shell with depth of edge beam to span ratio > 0.04 could be designed by beam theory.

Beam theory underestimates the vertical deformation very much. From Table 5.1 it is evident in this respect two theories can not be compared except for/few cases. Horizontal and rotational deformation have higher magnitudes than those of reinforced concrete shell.

Table 5.1 COMPARISON OF BEAM AND CLASSICAL THEORIES

L/R	$\emptyset_{\mathbf{k}}$	d/L e _O	e _m	e) - <u>-</u>		δ	Remark
	-	đ	d p	PSC	RCC	PSC 2	RCC
3 •33	30 40 45	0.03 0.4	0.4 0.4		5.5 27.75 45.82	-0.135 0.0715 0.105	0.89 0.388 0.262
5.00	30 40 45	0.03 0.4	0.4 0.4	30.00 23.60 15.00		0.078 0.14 0.16	1.030 0.800 0.590
10.00	30 40 45	0.03 0.4	0.4 0.4	69.60 69.2 72.0		0.6 0.524 0.458	1.08 1.156 1.43
3.33	30	0.05*1 0.04*1 0.03 0.02	0.4 0.4	56.5 50.0 31.0 12.2	9.1 4.0 5.55 11.5	0.016 -0.123 -0.135 -1.1	1.510 1.250 0.89 0.73
3 . 33	30	0.0 0.0 0.4 0.03 0.4 0.2 0.4	0.0 0.4 0.0 0.4 0.4 0.4 0.2*2	42.7 4.5 31.0 41.0 40.0		- - - - - - - - - - - - - - - - - - -	
3 •33	30	0.3 0.4	0.8 0.6 0.4 0.4 0.2 0.0	31.0 14.0	-	0.4 0.307 -0.135 -1.77	- - - 0.89

E = Maximum percentage of error in the variation of $\tau_{\mathbf{x}}$ by beam theory.

^{* 1} Tension occurs at the shell edge by classical theory

² At the bottom of edge beam theory gives compression while classical theory gives tension

³ Beam theory gives tension at crown.

5.2 EFFECT OF PRESTRESSING FORCE:

The reduction of stress resultants depends not only on the prestressing force but the combination of depth of the edge beam end and mid eccentricities. Difference between beam and classical theories decreases with sag. For shells, provided with axially prestressed edge beam, the beam theory agrees with classical theory within an error of 5 to 10 percent.

Prestressing force is the most influencial parameter for the disagreement between beam and classical theories.

With suitable combinations of depth of edge beam end and mid eccentricities. N_x and N_x decrease with the increase of prestressing force. Beam theory gives the same feature. In classical theory prestressing has a very sensitive influence in the region of edge beam for the distribution of N_x. Its influence in the distribution of N_x by classical theory, is not significant while beam theory indicates its significant influence. My changes its sign and Ny decreases with the increase of prestressing force by both the theories. Shells with the coefficients of prestressing force > 0.4 could not be designed by beam theory.

The downward vertical deformation for non prestressing shell gradually decreases and ultimately attains upword deformation with the increase of prestressing force for both the theories. Rotational and horizontal deformations

increased with prestressing.

Beam theory assumes loading is uniformly distributed over the surface or chord area of the shell and neglects the relative displacements within each transverse section. These are not correct for shells with edge beam. For reinforced concrete shell the corrective line loads for edge beam are not much while for prestressed shell it has high values. The vertical displacements, estimated by beam theory for reinforced concrete shell is well comparable with that by classical theory while for prestressed shell this is not true. Though beam theory assumes that the rotational and horizontal deformations are negligible, which are not correct, but this assumption does not have much influence on the agreement of beam and classical theories as the restraints against these deformations produces the stresses which are in a nature to compensate each other. The most influential point is the discrepancy in vertical deformation. This discripency increases with the corrective line loads due to edge beam, which inturn directly proportional to the prestressing force. This is the main reason of the disagreement between the shell and beam theory for reasonably higher values of prestressing force.

5.3 PROPOSED METHOD:

The proposed method which is detailed in 6-steps, given article 4.6, has been formulated considering the variation of external load with Sin $\frac{n\pi x}{n\pi x}$ in longitudinal direction. The arch analysis portion, developed by strain energy method, is simple and less tedious. The errors in the computations with the help of tables which are given in appendix are small. Maximum number of simultaneous equations to be solved is 3 which is not all difficult for compution. Shell with any type of boundary conditions can be analysed by this method. The values of the coefficients, given in the tables, are independent of structural geometry of the shell and can be used for any values of span to radius ratio, depth of edge beam to span ratio, thickness to radious ratio and semicentral angle. Stress resultants at any longitudinal section can be determined by this method.

5.4 SUMMARY OF THE CONCLUSIONS:

1. Reimforced concrete shell with any values of R/t, d/L and \emptyset_k can be analysed by beam theory with the proposed method for L/R more than 3.

Single barrel shells, exterior and interior of multibarrel shells could be analysed by beam theory with the proposed method.

2. The error varies 4 to 10% for $N_{\rm X}$, 2 to 12% for $N_{\rm X}$ %. Maximum error occurs for shells having higher values of L/R with lower values of $\emptyset_{\rm k}$ and lower values of $\emptyset_{\rm k}$ with higher values of L/R.

Errors can be minimised with the proper selection of \emptyset_k as higher the value of L/R higher the value of \emptyset_k .

- 3. Prestressing has very sensitive influence on $N_{\mathbf{x}}$ in the edge beam while in the shell portion it is not such sensitive. Change of $N_{\mathbf{x}}$ with the introduction of pressing is not much significant. Influence of prestressing on $N_{\mathbf{y}}$ and $M_{\mathbf{y}}$ has same level of influence throughout the shell.
- 4. The reduction of stress resultants is not dependent only on p but on the suitable combination of d/L, $e_{\rm o}/d$ and $e_{\rm m}/d$. along with p .

Shells with axially prestressed edge beam can be analysed by the proposed method of beam theory. Shells provided with $d/L \gg 0.04$ and coefficient of prestressing force more than or equal 0.4 can not be analysed by beam theory.

6. The proposed method to analyse shell with beam theory is systematic, simple and less tedious. Tables, which are given to analyse a shell can be used for any values of

R/L and R/t and \emptyset_k . The stress resultant at any longitudinal section can be estimated by this method.

5.5 SCOPE OF FURTHER WORK:

- 1. Analysis of cylindrical shells by beam theory with different support and loading condition.
- 2. It appears worthwhile to extend the analysis for the outer shells for 6 boundary conditions for beam theory and eight boundary conditions for classical theory and to study the effect of higher Fourier terms on the proposed method.
- 3. Development of the algorithm of to attain the optimal design criterion of cylindrical shells by beam theory.
- 4. It is worth-while to prepare the design table of prestressed concrete cylindrical shell by beam theory.

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APPENDIX

A.1 SHELL COEEFICIENTS

Table A.1.1 SHELL COEFFICIENTS

 $F = 2\overline{R} \begin{cases} (\emptyset \text{ Measured from the Crown}) \\ (aB_1 - bB_2) \cos \beta_1 \emptyset \cosh \alpha_1 \emptyset - (aB_2 + bB_1) \sin \beta_1 \emptyset \sinh \alpha_1 \emptyset \\ (cB_3 - dB_4) \cos \beta_1 \emptyset \cosh \alpha_1 \emptyset - (cB_4 + dB_3) \sin \beta_1 \emptyset \sinh \alpha_1 \emptyset \end{cases}$

						
F	R	^B 1	^B 2	B ₃	B ₄	
$^{ m M}$ ø	$-\frac{2D}{r^2}$ sin kx	$n^2(1+V)$	n ²	n ² (V-1)	n ²	
$^{ m M}_{ m x}$	2Dk ² sin kx	. 1	0	1	0	
$Q_{\mathbf{X}}$	$-\frac{2Dk}{}^{3}\cos kx$		1	-1	1	
$^{ m N}$ ø	$\frac{4Drk^4}{\sqrt{2}}$ sin kx	0	1	0	-1	
$^{\mathbf{x}}$	$-\frac{4Drk^4}{\sqrt{3}}$ sin kx	-1	1 + √	1	1 - 1	
$Q_{\mathbf{X}}^{1}$	$-\frac{2Dk^3}{}\cos kx$	V + 2	2	V - 2	2	
u _x	4Drk ³ cos kx	- 1	1 + √		+ - 1	
wr	2 sin kx		0		0	

with $\beta = 0$, so that $D = Et^3/12$

$$k = \underbrace{n\pi.}_{1}$$

Table A.1.2 SHELL COEFFICIENTS (O Measured from the Crown)

	8	70
	cosh α_1	2 K+0
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•	~	0-(cB4+dB3) sin 810
:	sinh $\alpha_1 \beta - (aB_2 + bB_1)$ sin $\beta_1 \beta$	sinh $\alpha_1^{i} \emptyset$
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	${ m B_4}$	2 2m	$m_2 - n_2 (1 + V)$	C7 E1	-n2-m2(1+V)	$\frac{2\beta_{4}}{r} + \frac{(\bar{R} \cdot B_{2})_{v}}{r} \frac{2\alpha_{1}'}{r} + \frac{(\bar{R} \cdot B_{3})_{v}}{r} \frac{2\beta_{1}'}{r} \frac{(\bar{R} \cdot B_{4})_{v}}{r}$	8	
	$^{\mathrm{B}_{3}}$	-(m ₂ +n ₂)	-m ₂ (1+1)-n ₂	n2	$-m_2+n_2(1+V)$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	βαί	
the cost of the standard of the cost of th	\mathtt{B}_2	m ₁ + m ₁	m ₁ t n ₁ (1-V)	, a	n,-m, (1-V)		β	
or alm mire alo	Fg.	tu-tu	m ₁ (1-V)-n ₁	u u	m ₁ + n ₁ (1-V)	$\frac{2\alpha_1}{r} + \frac{(R \cdot B_1)_{V}}{r}$	8	= Eh ² /12
1 200 / Pan-Earl	R	$-\frac{Dk^3}{(V)^3/2}$ gin kx	$-\frac{20k^3}{(\sqrt{)}^{3/2}}\sin kx$	$\frac{4Drk^4}{(V)^5/2}\cos kx$	$\frac{4Dr^k}{Eh.(\sqrt{V})^7/2}$	sin kx	2Dk cos kx	$\vec{y} = 0$, so that $D = Eh^3/12$
	É	Øo	-%	$N_{X}\emptyset$	A	-0	M_{∞} M	

1 Observe that one part of the coefficients B, etc., for are obtained from v, that is

$$(\bar{R} \cdot B_1)_{V}/r = \frac{4Dk^3}{Eh(V)^7/2}n_1 + n_1 (1 - V)$$

The subscript v indicates uy

1.2 COEFFICIENTS OF THE TRANSVERSE SECTIONAL PROPERTIES

Table A.2.1 COEFFICIENTS FOR C.G.

(in	Ø degree)	ø (in rad)	Cos Ø	Sin Ø	k _z
	45.00	0.78540	0.70711	0.70711	0.09968
	42.50	0.74176	0.73728	0.67559	0.08921
	40.00	0.69813	0.76604	0.64279	0.07927
	37.50	0.65450	0.79335	0.60876	0.06988
	35.00	0.61087	0.81915	0.57358	0.06104
	32.50	0.56723	0.84339	0.53730	0.05277
	30.00	0.52360	0.86603	0.50000	0.04507
	27.50	0.47997	0.88701	0.46175	0.03795
	25.00	0.43633	0.90631	0.42262	0.03143
	22.50	0.39270	0.92388	0.38268	0.02550
	20.00	0.34907	0.93969	0.34202	0.02018
	17.50	0.30543	0.95372	0.30071	0.01548
	15.00	0.26180	0.96593	0.25882	0.01138
	12.50	0.21817	0.97630	0.21644	0.00791
	10.00	0.17453	0.98481	0.17365	0.00507
	7.50	0.13090	0.99144	0.13053	0.00285
	5.00	0.08727	0.99619	0.08716	0.00127
	2 50	0.04363	0.99905	0.04362	0.00032
	-0.00	-0.00000	1.00000	-0.00000	-0.00000

Table A.2.2 <u>COEFFICIENTS FOR MOMENT OF INERTIA</u> (Ø Measured from crown)

ø _k * 	kΙ		kĮ		^k I	$ \emptyset_{\mathrm{k}}^{*} $	k _I €
45.00	0.01216	44.50	0.01152	44.00	0.01091	43.50	0.01032
43.00	0.00976	42.50	0,00922	42.00	0.00871	41.50	0.00822
41.00	0.00775	40.50	0.00730	40.00	0.00687	39.50	0.00647
39.00	0.00608	38.50	0.00571	38,00	0.00536	37.50	0.00502
37.00	0.00470	36.50	0.00440	36.00	0.00411	36.50	0.00334
35.00	0.00358	34.50	0.00334	34.00	0.00311	33.50	0.00289
33.00	0.00269	32.50	0.00249	32.00	0.00231	31.50	0.00214
31.00	0.00198	30.50	0.00182	30.00	0.00168	29.50	0.00155

^{*} p_k , given in degrees

Table A.3.1 COEFFICIENTS FOR ARCH GALCULATIONS

ϕ_{K}	kmv	Kmh	k mm	^k mvħ	Kmvm	kmhm	$^{ m K}$ nv	k ⁿ th
45.00	12.84956	0.84956	18,84956	-2.57522	13,32865	-3.64191	0.28540	1.28540
44.50	12,47837	0.80552	18,64012	-2:47193	13.06503	-3.42675	0.27675	1.27680
44.00	12.11272	0.76325	18.43068	-2.37148	12,83302	-3.41398	0.26825	1,26764
43.50	11.75265	a. 72272	13.22124	-2.27384	12,54267	-3.30329	0.25990	1.25853
43.00	11,39819	0.68386	18,01180	-2,17896	12,28402	-3.19497	0.25171	1.24927
42.50	11.04940	0.64664	17.80236	-2,08682	12.02710	-3.08889	0.24367	0.123986
42.00	10.70631	0.61100	17.59292	-1.99739	11.77196	-2.98505	0.23478	1.23030
41.50	10,36895	0.57690	17.38348	-1.91062	11,51864	-2,88342	0.22804	1,22058
41.00	10.03735	0.54430	17.17404	-1:82647	11.26718	-2,78400	0,22045	1.21072
40.50	9.71155	0.51313	16.96460	-1:74492	11.01763	-2.68677	0,21301	1.20070
40.00	9.39156	0.48337	16.75516	-1.66592	10,77001	-2.59170	0,20573	1.19054
39.50	9.07742	0.45497	16.54572	-1.58943	10,52437	-2.49879	0.19859	1.18022
39.00	8,76914	0.42788	16,33628	-1.51541	10:28075	-2,40801	0.19160	1.16975
38.50	8,46675	0.40206	16,12684	-1:44383	10.03919	-2:31935	0.18477	1.15914
38.00	8,17024	0.37746	15.91840	-1.37464	9.79973	-2:23279	0.17808	1.14837
37.50	7.87965	0.35460	15.70796	-1.30781	9.56240	-2,14831	0.17154	1.13746

Contd....

ø.	$k_{\mathbf{m}\mathbf{v}}$	^K mh	$k_{ m rmn}$	$k_{ m mvh}$	$k_{ m mvm}$	k _m hm	k_{nv}	k _n h
37.00	7.59497	0.33180	15.49852	-1.24328	9.32724	-2,06582	0.16514	1.12640
36.50	7,31621	0.31065	15,28908	-1.18103	9.09429	-1.98551	0.15889	1.11520
36.00	7.04337	0.29056	15.07964	-1.12100	8.86359	-1.90716	0.15279	1.10385
35.50	6,77646	0.27150	14,87020	-1.06315	8.63517	-1.83081	0,14683	1,09235
35.00	6.51547	0.25344	14.66076	-1.00745	8,40907	-1.75644	0.14102	1.03071
34.50	6,26040	0.23632	14.45132	-0.95385	8,18532	-1.68403	0.13535	1.06893
34.00	6.01123	0.22013	14.24189	-0.90230	7.96396	-1.61357	0.12982	1.05700
33.50	5.76796	0,20482	14.03245	-0.85276	7.74502	-1.54503	0.12443	1.04494
33.00	5.53057	0.19035	13.82301	-0.80519	7.52855	-1.47839	0.11919	1.03273
32.50	5.29905	0.17670	13.71357	-0.79554	7.31456	-1.41365	0.11408	1.02039
32.00	5.07337	0,16383	13.40413	-0.71577	7.10310	-1.35072	0.10911	1.00790
31.50	4.85352	0.15170	13.19469	-0.67384	6.89420	-1.28964	0.10428	0.99528
31.00	4.63946	0.14030	12,98525	-0.63369	6.68790	-1.23038	0.09958	0.98253
30.50	4.43117	0.12957	12,77581	-0.59530	6.48421	-1:17291	0:09503	0.96964
30.00	4.22862	0:11951	12.56637	-0.55860	6,28381	-1.11720	0.09059	0.95661
29,50	4.03178	0.11006	12.35693	-0.52356	6.08484	-1,06324	0.08629	0.94346

A. 4 COEFFICIENTS FOR STRESS RESULTANTS

Table A.4.1 COEFFICIENTS FOR STRESS RESULTANTS (No)

$m{arphi}_{\mathbf{k}}$	c _{vn}	c _{hn}	$\emptyset_{\mathbf{k}}$	c _{vn}	c _{hn}
45.00	-0.70711	-0.70711	42.50	-0.67559	-0.73728
40.00	-0.64279	-0.76604	37.50	-0.60876	-0.79335
35.00	-0.57358	-0.81915	32.50	-0.53730	-0.84339
30.00	-0.50000	-0.86602	27.50	-0.46175	-0.88701
25.00	-0.42262	-0.90631	22,50	-0.38268	-0.92388
20.00	-0.34202	-0.93969	17.50	-0.30071	-0.95372
15.00	-0.25882	-0.96593	12.50	-0.21644	-0.97630
10.00	-0.17365	-0.98481	7.50	-0.13053	-0.99144
5.00	-0.08716	-0.99619	2.50	-0.04362	-0.99905
0.00	-0.00000	-1.00000			

Table A.4.2 COEFFICIENTS OF STRESS RESULTANTS (Mg)

ø	c _{vm}	$\mathbf{c}_{ ext{hm}}$	Ø	$\mathbf{c}_{\mathtt{Vm}}$	c _{hm}
		Ø _I	=32.50		
30: 00	0.00000	0.00000	32.50	0.00000	0.00000
25.00	0.07738	-0.04028	30.00	0.03730	-0.02263
20.00	0.15798	-0.07367	25.00	0.11468	-0.06292
15.00	0.24118	-0.09990	20.00	0.19528	-0.09630
10.00	0.32635	-0.11878	15.00	0.27848	-0.12253
5.00	0.41284	-0.13017	10.00	0.36365	-0.14142
0.00	0.50000	-0.13397	5.00	0.45140	-0.194280
			0.00	0.53730	-0.15661

Contd.....

Ø	c^n	$\mathbf{c}_{\mathtt{hm}}$	Ø	cvm	chm
$\emptyset_{k} = 35.00$			Ø _k =37.50		
35.00	0.0000	0,00000	37.50	0.0000	0.0.0000
30.00	0.73580	-0.04687	35.00	0.0351	8 -0.02580
25.00	0.15096	-0.08716	30.00	0.1027	6 -0.07267
20.00	0.23156	-0.12054	25.00	0.1861	4 -0.11295
15.00	0.31476	-0.14677	20.00	0.2667	4 -0.14634
10.00	0.39993	-0.16566	15.00	0.3499	4 -0.17227
5.00	0.48642	-0.17704	10.00	0.4351	1 -0.19145
0.00	0.57358	-0.18035	5.00	0.5216	1 -0.20284
			0.00	0.6087	6 -0.20665
$g_{k} = 40.00$			Ø _k =42.50		
40.00	0.00000	0.0000	42.50	0.0000	0,00000
35.00	0.06921	-0.05311	40.00	0.0328	0.02877
30.00	0.14279	-0.09998	35.00	0.1020	1 -0.08187
25.00	0.22017	-0.14026	30.00	0.1755	9 -0.12875
20.00	0.30077	-0.17365	25.00	0.2529	7 -0.16903
15.00	0.38397	-0.19988	20.00	0.3335	7 -0.20242
10.00	0.46914	-0.21876	15.00	0.4167	7 -0.22865
5.00	0.55563	-0.23015	10.00	0.5019	4 -0.24753
0.00	0.64279	-0.23396	5.00	0.5884	3 -0.25892
			0.00	0.6755	9 -0.26272
		$ ot\!\!\!/ p_1$	=45.00		
45.00	0.0000	0.00000	20.00	0.3650	9 -0.23259
40.00	0.06432	-0.05894	15.00	0.4482	9 -0.25882
35.00	0.13353	-0.11205	10.00	0.5334	6 -0.27770
30.00	0.20711	-0.15892	5.00	0.6199	5 -0.28909
25.00	0.28449	-0.19920	0.00	0.7071	1 -0.29289